

The r-process and electromagnetic emission from neutron star mergers

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with

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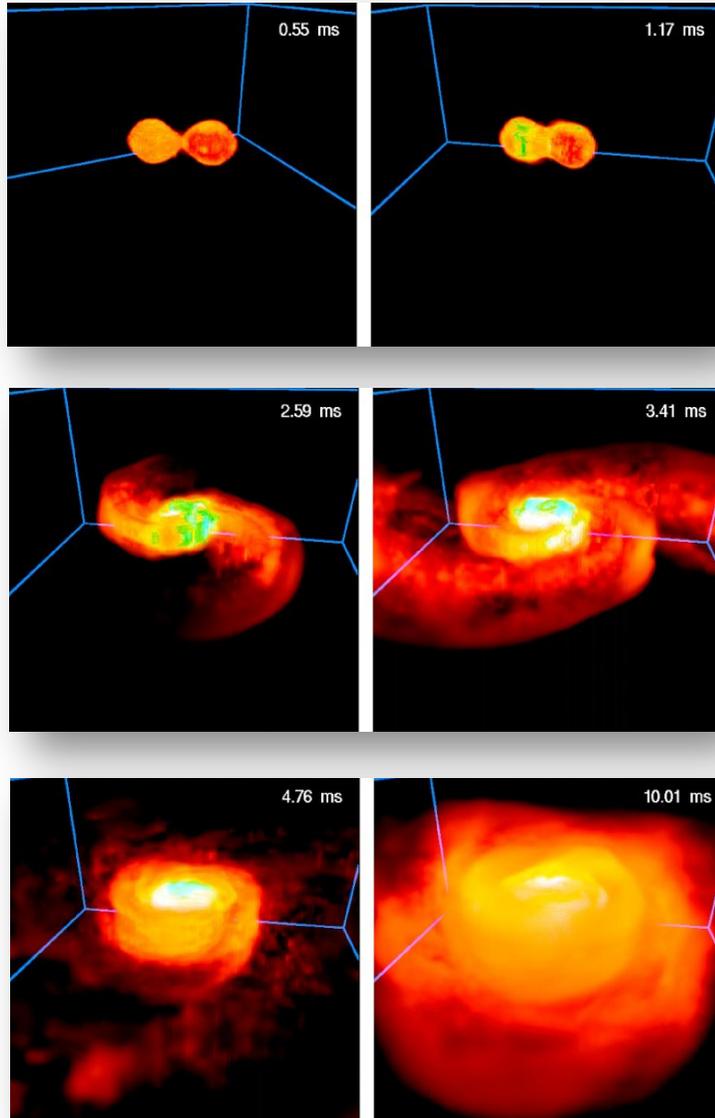
origin of gold (r-process elements) is still unknown...



www.cartier.jp

NS mergers: sRGB and r-process

www.mpa-garching.mpg.de



- ❖ coalescence of binary NSs
expected $\sim 10 - 100$ per Myr in the Galaxy (possible sources of short GRB as well as r-process)
- ❖ first ~ 0.1 seconds
dynamical ejection of n-rich matter up to $M_{\text{ej}} \sim 10^{-2} M_{\odot}$
- ❖ next ~ 1 second
neutrino or viscously driven wind from the BH accretion torus up to $M_{\text{ej}} \sim 10^{-2} M_{\odot} ??$

NS mergers: sGRB and r-process

LETTER

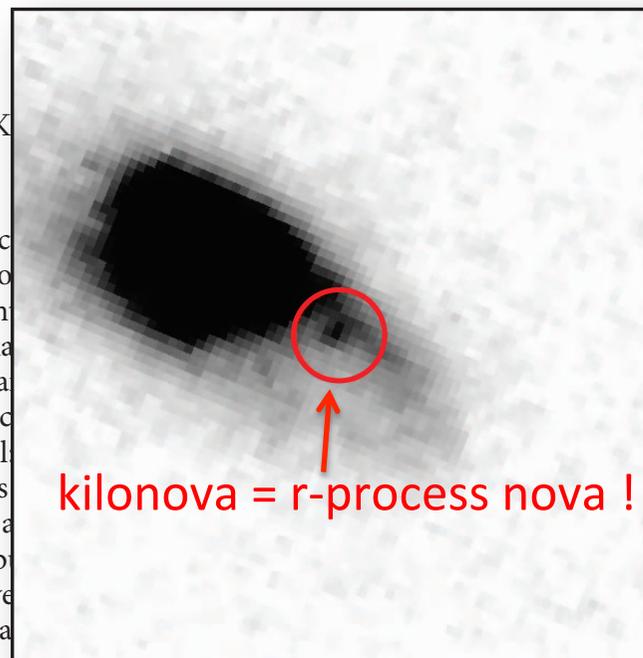
doi:10.1038/nature12505

A 'kilonova' associated with the short-duration γ -ray burst GRB 130603B

N. R. Tanvir¹, A. J. Levan², A. S. Fruchter³, J. Hjorth⁴, R. A. Hounsell³, K.

Short-duration γ -ray bursts are intense flashes of cosmic γ -rays, lasting less than about two seconds, whose origin is unclear^{1,2}. The favoured hypothesis is that they are produced by a relativistic jet created by the merger of two compact stellar objects (specifically two neutron stars or a neutron star and a black hole). This is supported by indirect evidence such as the properties of their host galaxies³, but unambiguous confirmation of the model is still lacking. Mergers of this kind are also expected to create significant quantities of neutron-rich radioactive species^{4,5}, whose decay should result in a faint transient, known as a 'kilonova', in the days following the burst⁶⁻⁸. Indeed, it is speculated that this mechanism may be the predominant source of stable r-process elements in the Universe^{5,9}.

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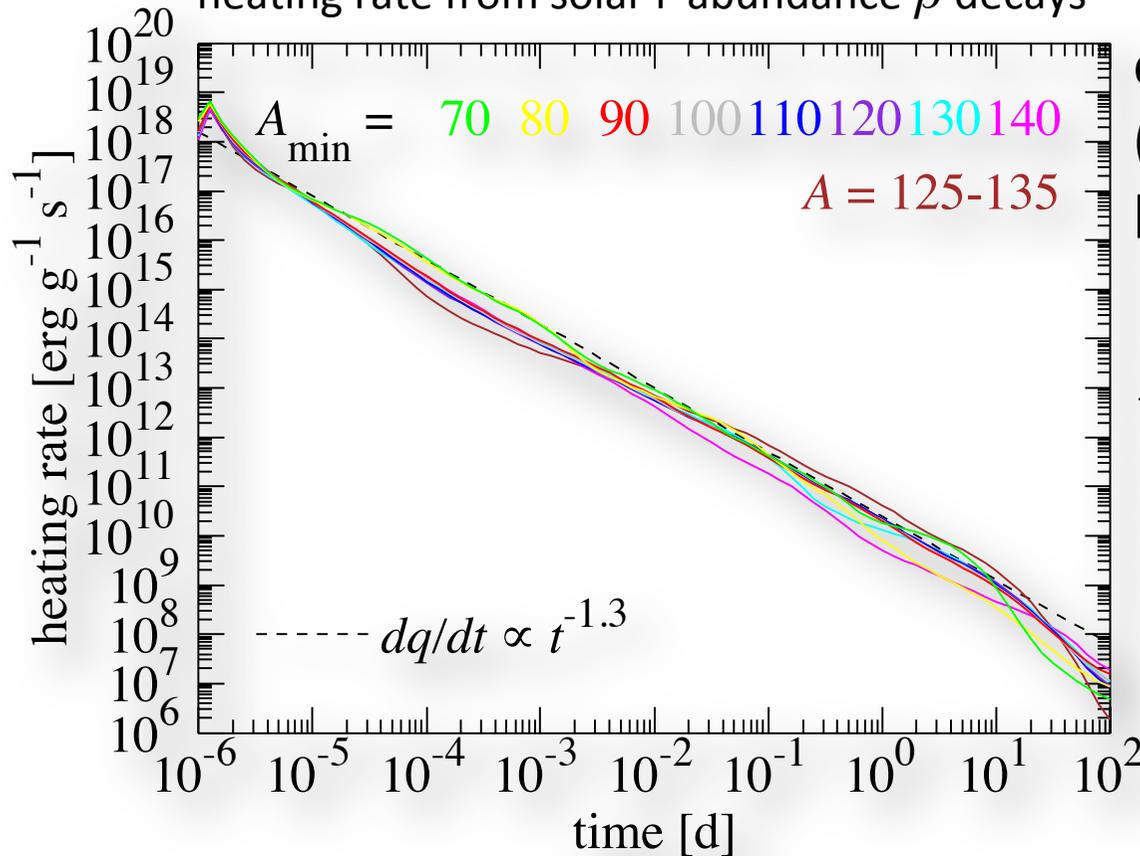


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Tanvir+2013, Nature, Aug. 29

r-process novae (kilonovae)

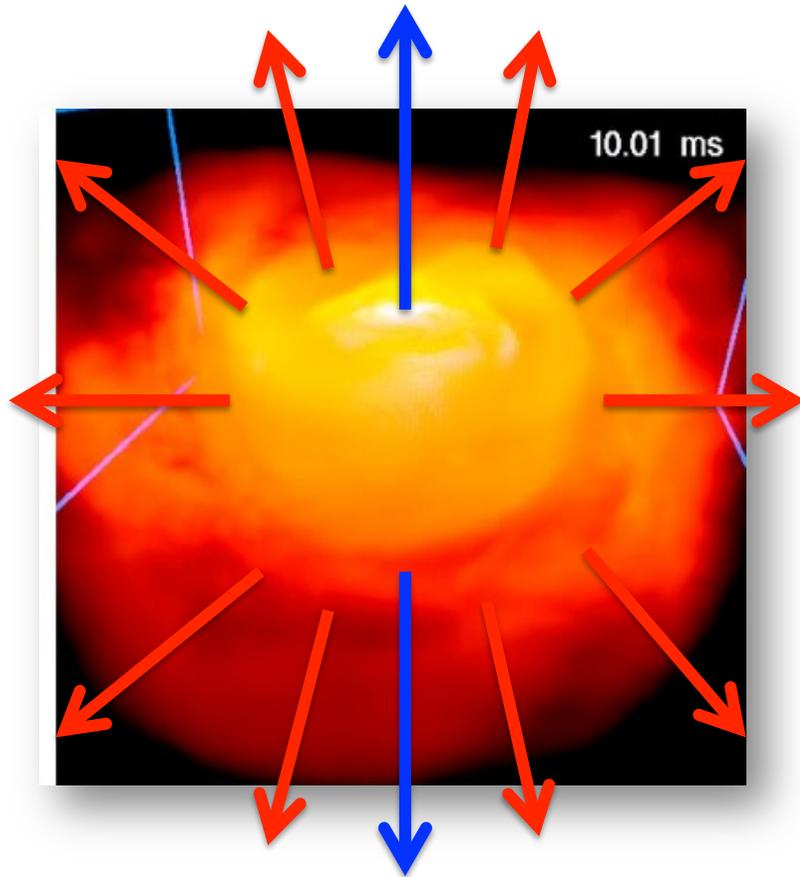
Wanajo+2014;
heating rate from solar r-abundance β -decays



electro-magnetic transients
(macronova, Kulkarni 2005;
kilonova, Metzger+2010)

- ❖ heating from decays of radionuclides well scaled as $dq/dt \sim t^{-1.3}$

EM counterparts of GW signals



GW signal can be spatially resolved only $\sim 100 \text{ deg}^2$ by KAGRA/a.LIGO/a.Virgo (from 2017)

→ EM counterparts are needed

❖ SGRBs

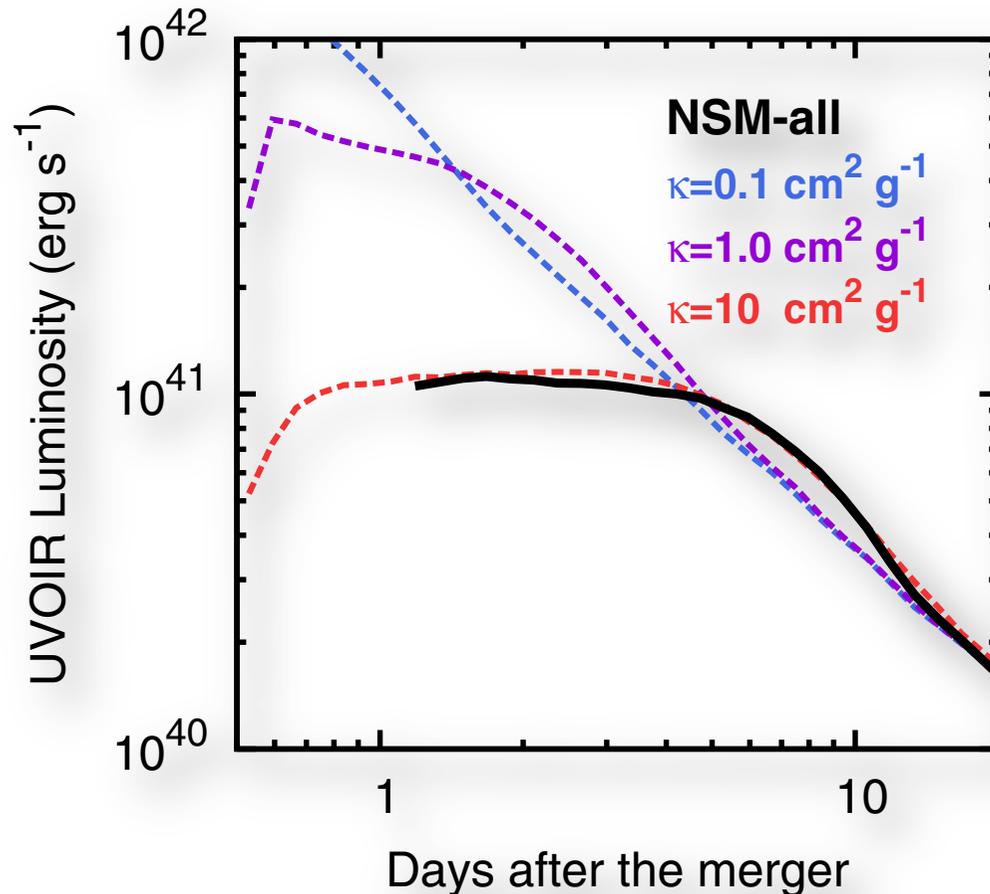
events should be restricted due to narrow beaming

❖ r-process novae

detectable (by, e.g., Subaru/HSC) from all directions!

lanthanide curtain for r-process novae

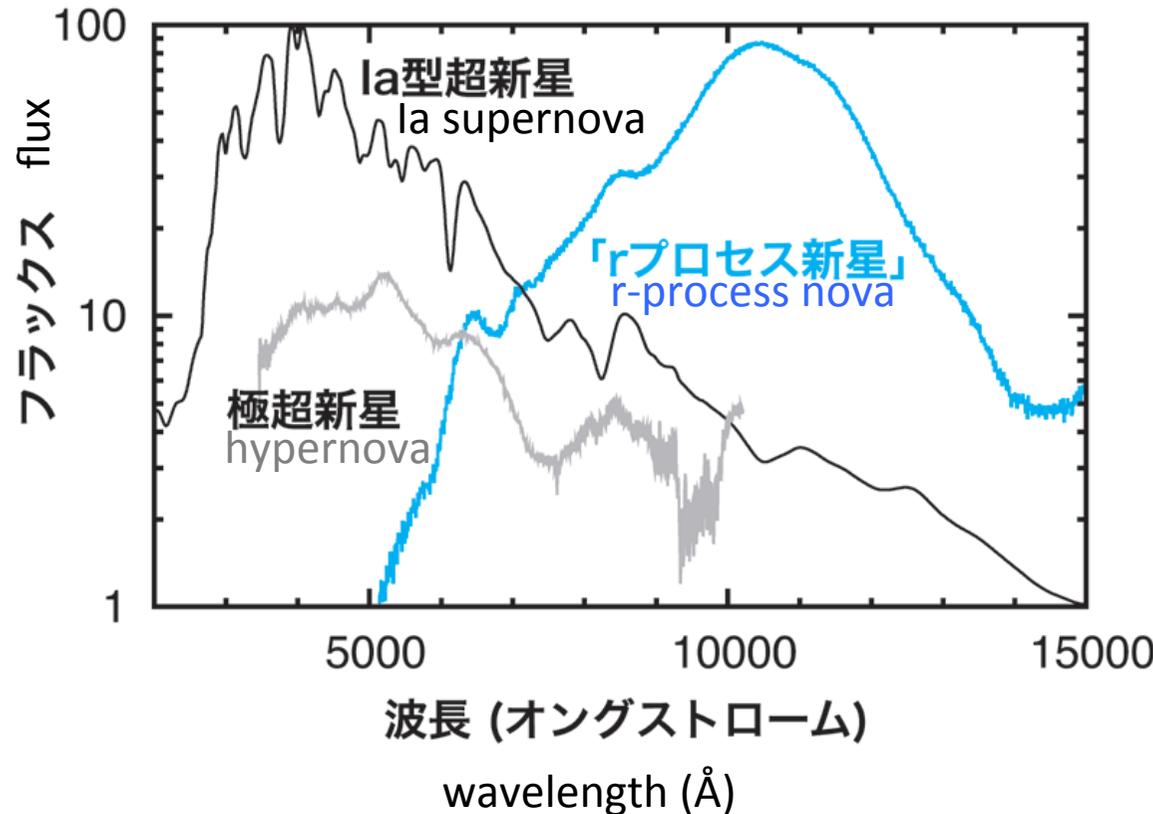
Tanaka & Hotokezaka 2013;
NS+NS models



- ❖ large opacities of lanthanides ($A > 130$);
~ 100 times greater than those of iron group elements
- ❖ brightness of r-process novae should be highly dependent on the nucleosynthetic abundances

what is a smoking gun of the r-process?

田中, 天文月報2014年1月



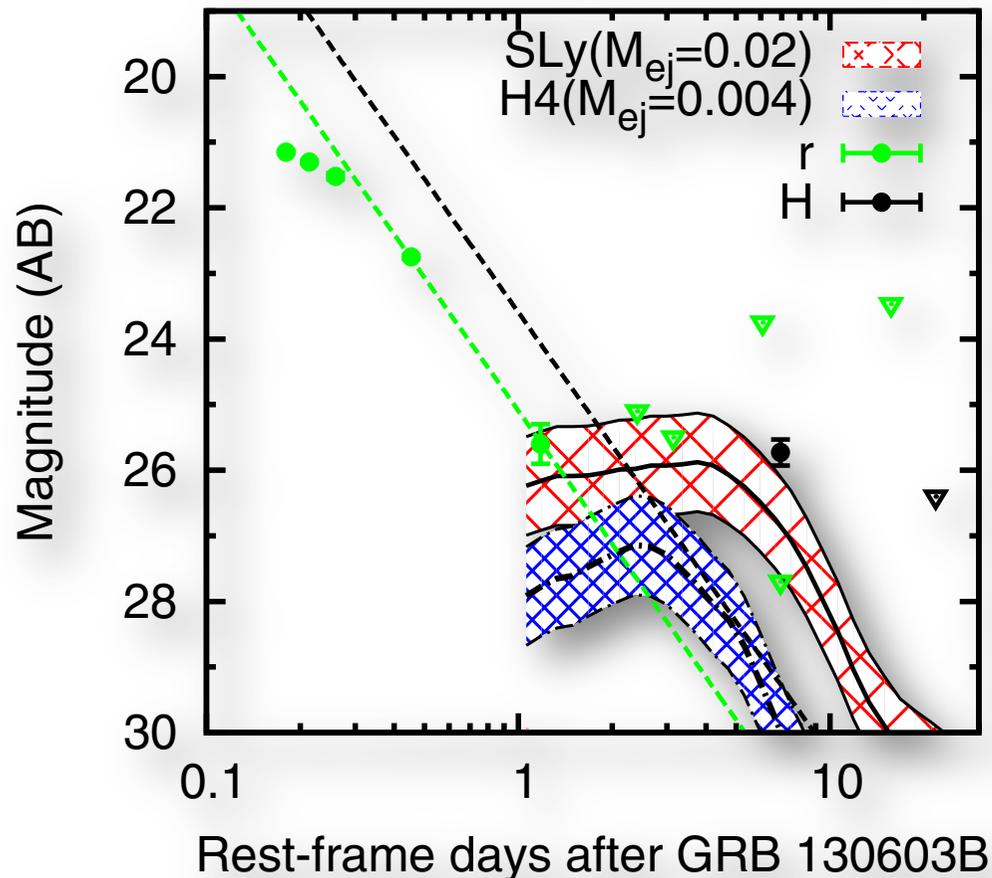
can we see r-abundances in the spectra?

❖ almost featureless because of too many bound-bound lines and Doppler shifts ($v/c \sim 0.1-0.3$)

❖ identification of red, featureless spectral shape can be an unambiguous evidence of an r-process

r-process nova in the SGRB afterglow?

Hotokezaka+Tanaka...+Wanajo 2013;
NS+NS models

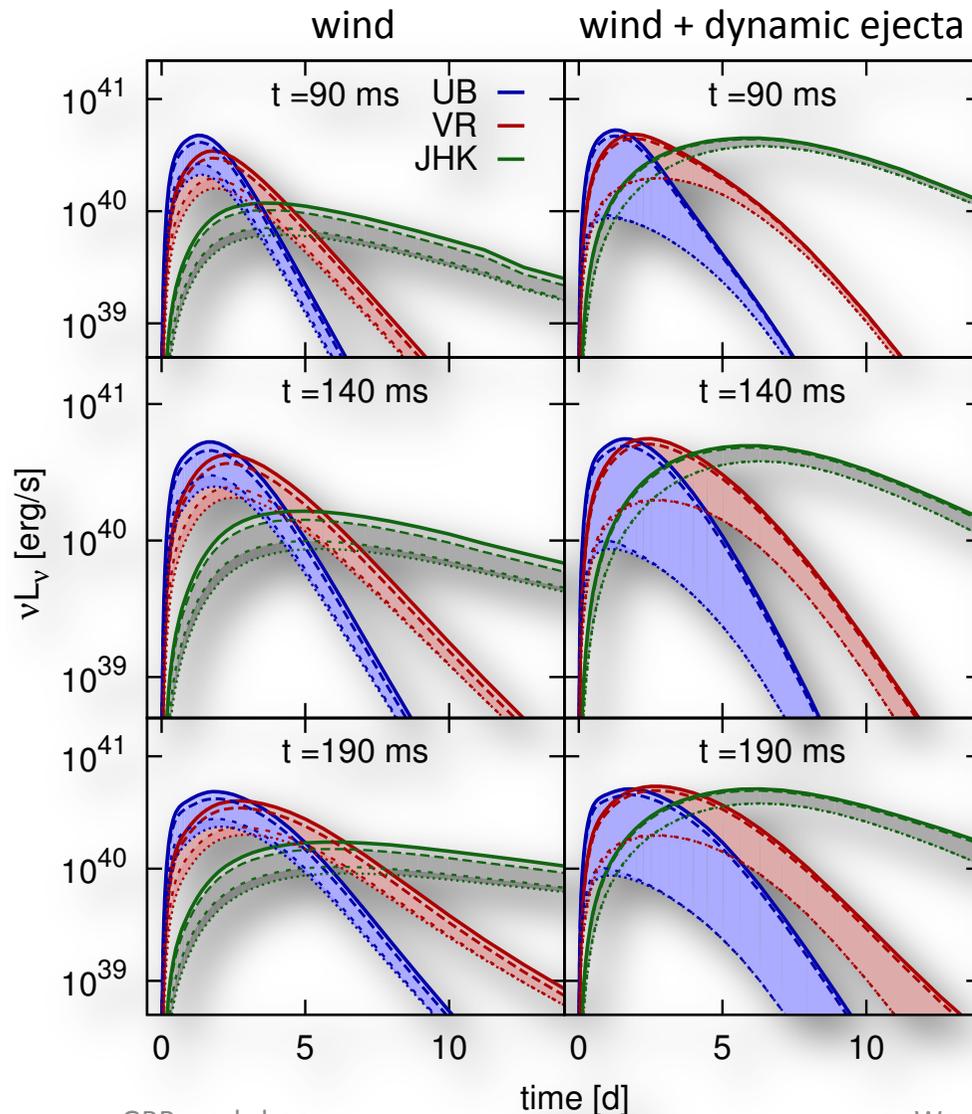


❖ late-time excess NIR flux
requires an additional
component (most likely an
r-process nova)

❖ the excess NIR indicates the
NS-NS ejecta with
 $M_{ej} \sim 0.02 M_{\odot}$

late-time wind and viewing angle

Martin+2015; kilonova light curves (0-90 deg)

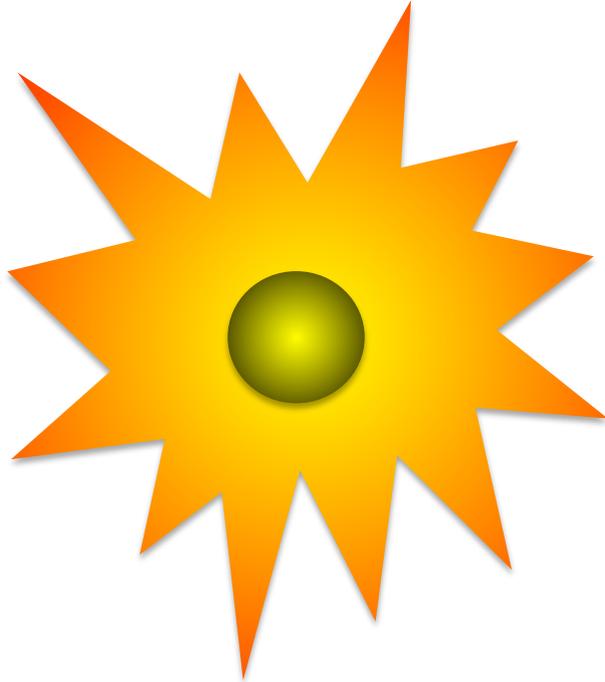


❖ effects of wind components as well as viewing angles (e.g., Martin+2015 with Newtonian simulations)

❖ wind component ($\sim 0-60$ deg) is lanthanide free ($A < 130$)

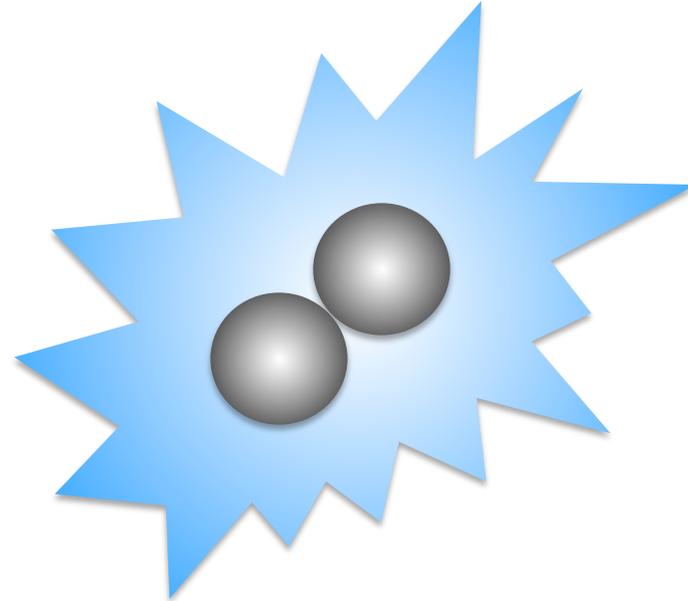
❖ r-process novae can be bright in UV and VR in the first few days; detailed nucleosynthetic information (abundances as function of viewing angle) should be important

where do we have neutrons?



core-collapse supernovae
(since Burbidge+1957;
Cameron 1957)

- ❖ n-rich ejecta nearby proto-NS
- ❖ not promising according to recent studies

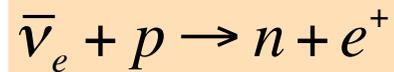


neutron-star mergers
(since Lattimer+1974;
Symbalisty+1982)

- ❖ n-rich ejecta from coalescing NS-NS or BH-NS
- ❖ few nucleosynthesis studies

SN neutrino wind: not so neutron-rich

❖ Y_e is determined by



❖ equilibrium value is

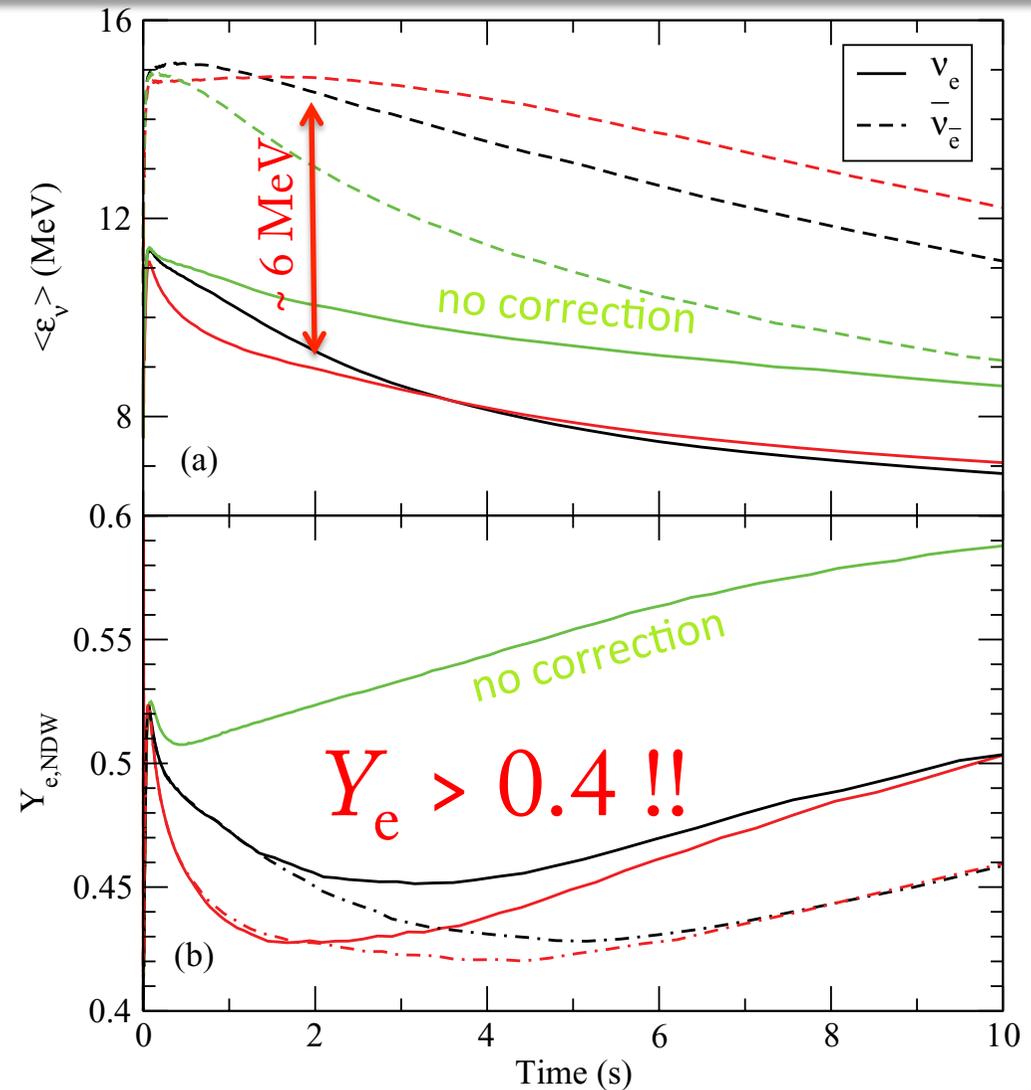
$$Y_e \sim \left[1 + \frac{L_{\bar{\nu}_e} \varepsilon_{\bar{\nu}_e} - 2\Delta}{L_{\nu_e} \varepsilon_{\nu_e} + 2\Delta} \right]^{-1},$$

$$\Delta = M_n - M_p \approx 1.29 \text{ MeV}$$

❖ for $Y_e < 0.5$ (i.e., n-rich)

$$\varepsilon_{\bar{\nu}_e} - \varepsilon_{\nu_e} > 4\Delta \sim 5 \text{ MeV}$$

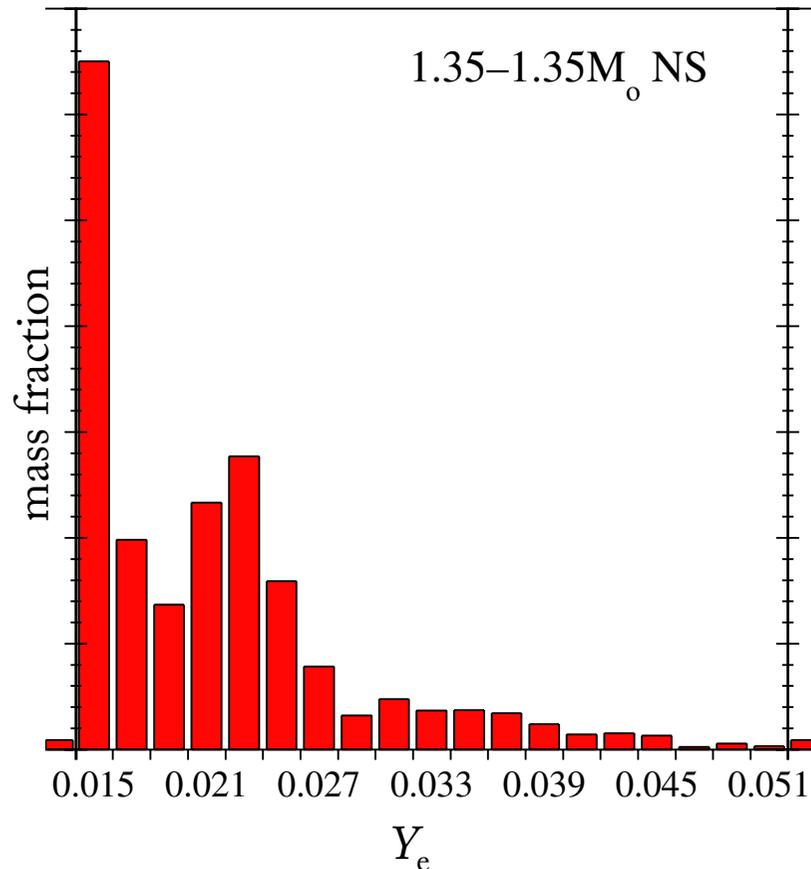
$$\text{if } L_{\bar{\nu}_e} \approx L_{\nu_e}$$



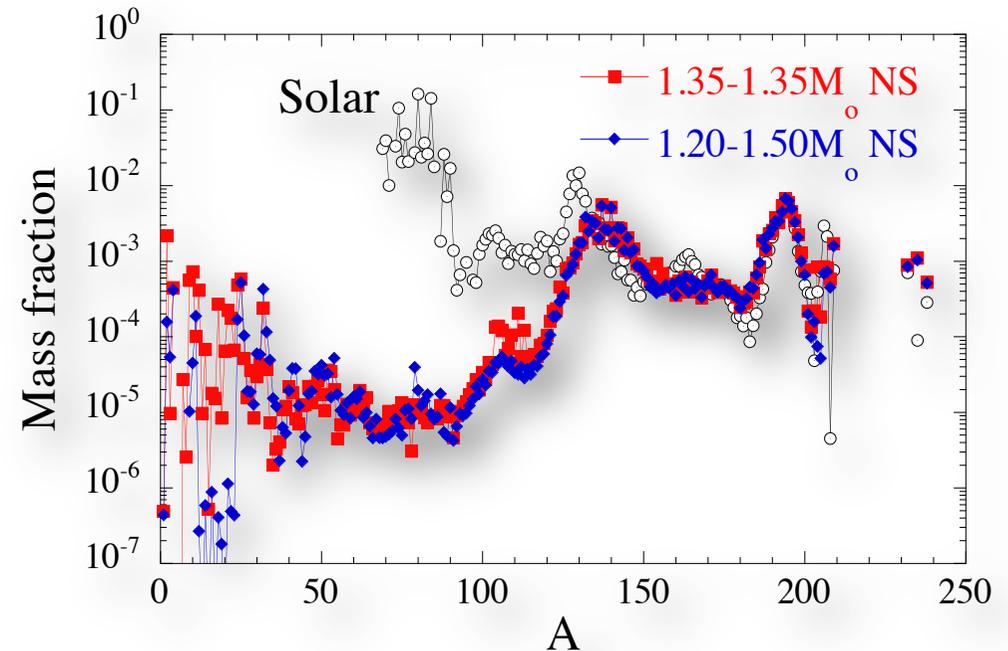
Roberts+2012

previous works: too neutron-rich ?

Goriely+2011 (also similar results by Korobkin+2011; Rosswog+2013)



tidal (or weakly shocked) ejection
of “pure” n-matter with $Y_e < 0.1$

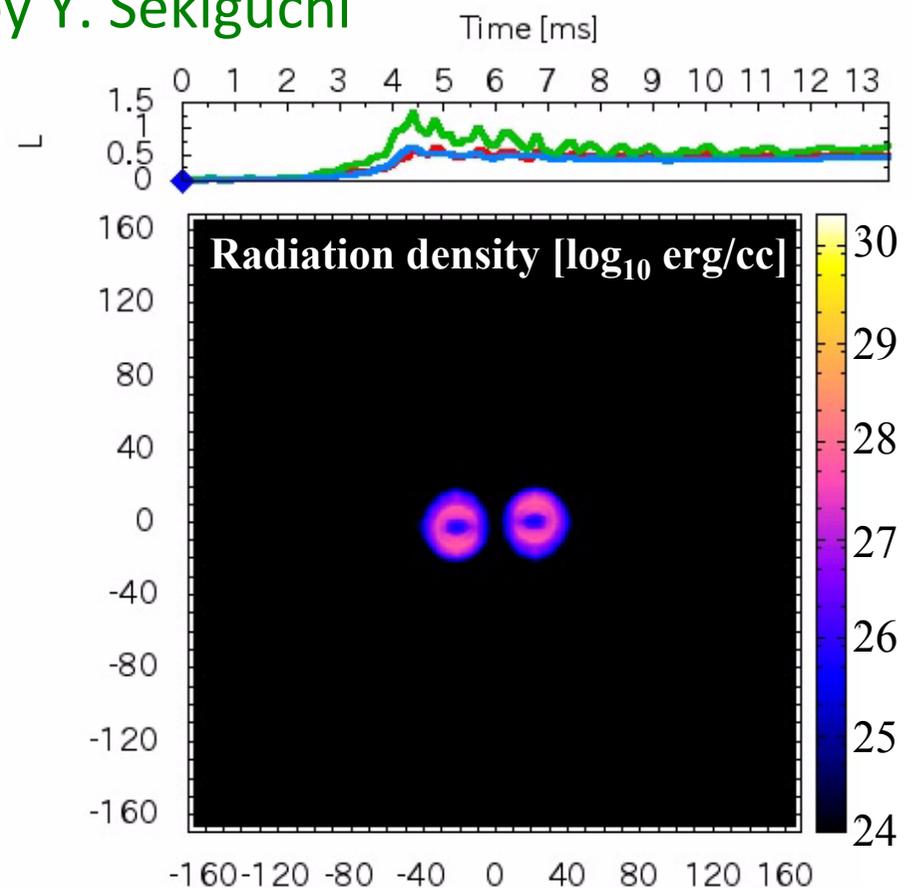
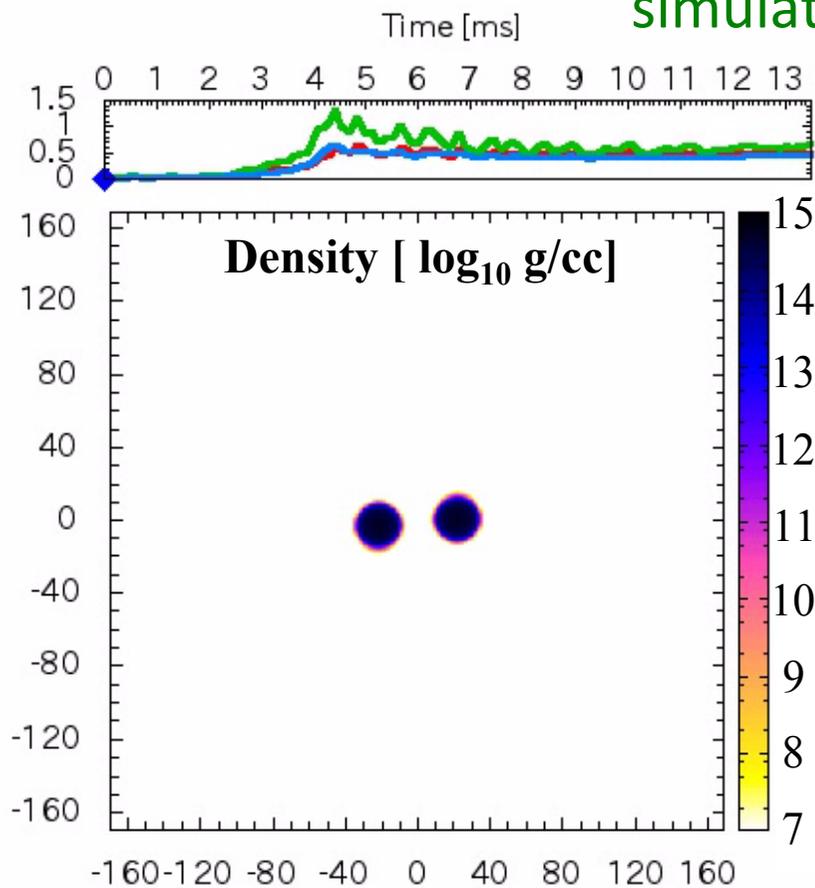


- ❖ strong r-process leading to fission recycling
- ❖ severe problem: only $A > 120$; another source is needed for the lighter counterpart

first simulation with full-GR and ν

- ▶ Approximate solution by Thorne's Moment scheme with a closure relation
- ▶ Leakage + Neutrino heating (absorption on proton/neutron) included

simulation by Y. Sekiguchi

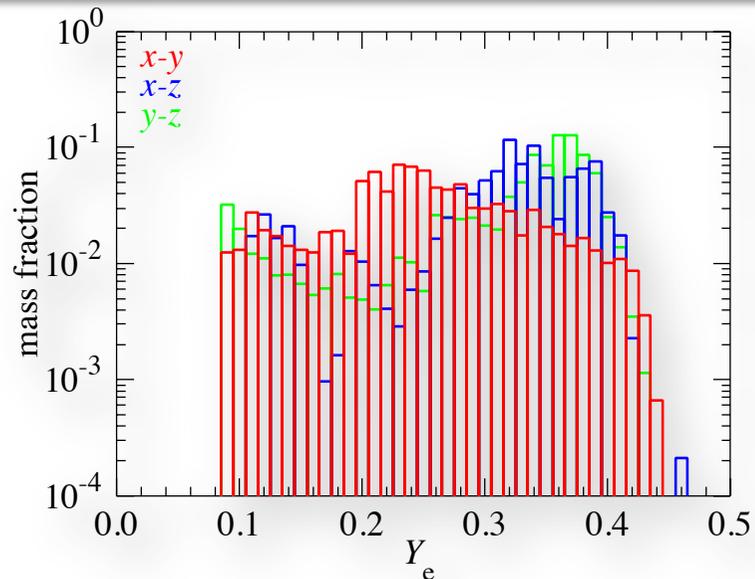


1.3+1.3 M_{\odot} neutron star merger with full-GR and neutrino transport (SFHo)

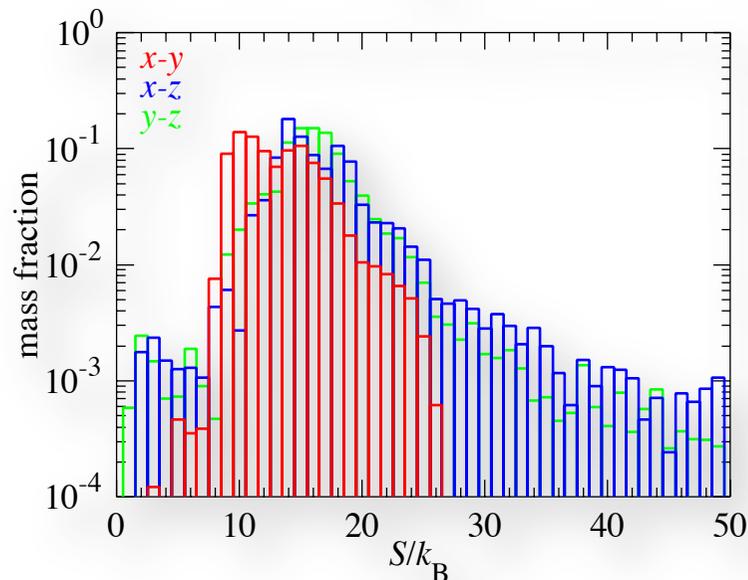
simulation by Yuichiro Sekiguchi



nucleosynthesis in the NS ejecta

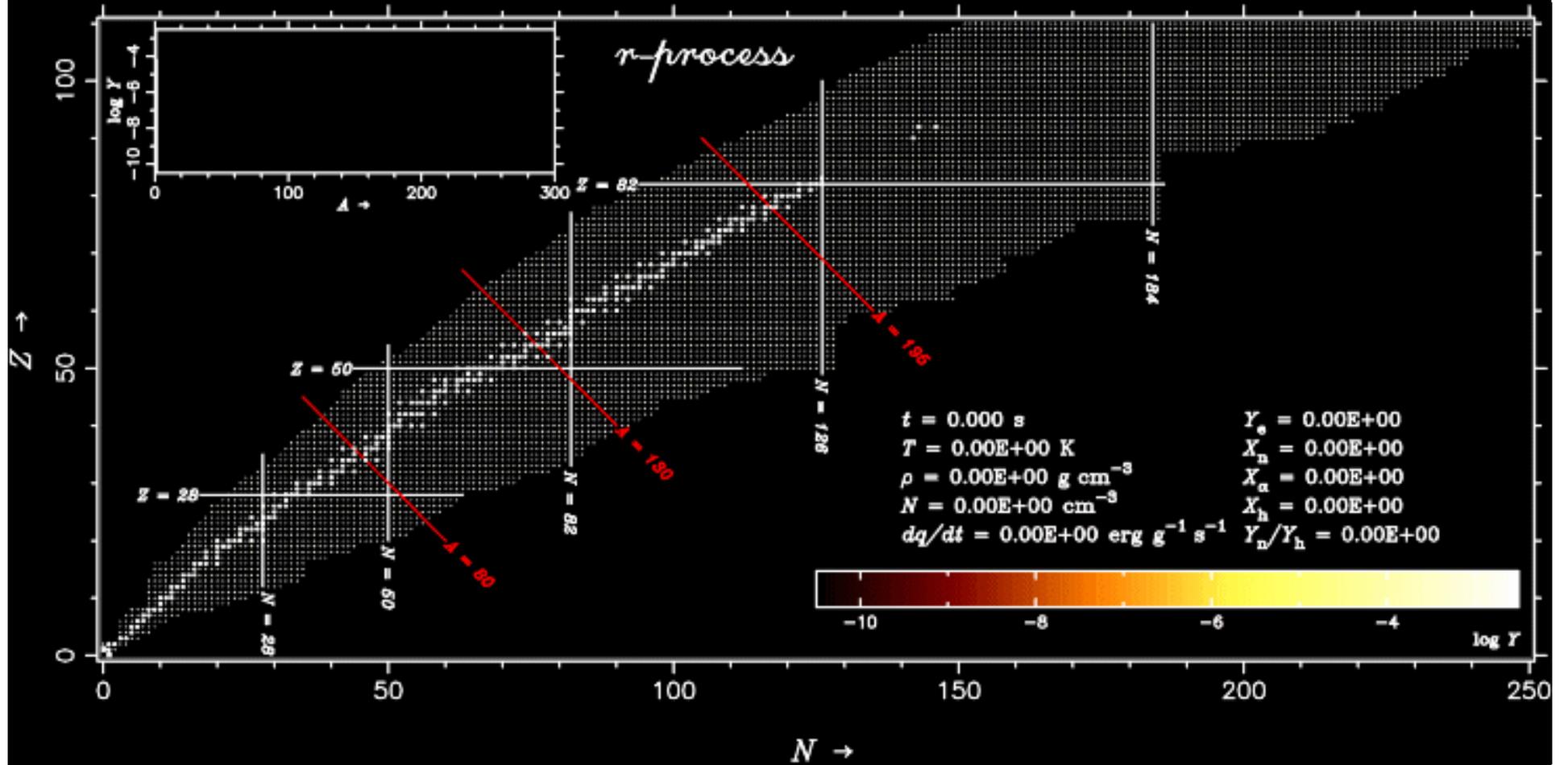


- ❖ higher and wider range of Y_e ($\sim 0.1-0.5$) in contrast to previous cases Y_e ($= 0.01-0.05$)
- ❖ values do not fully asymptote to $Y_e \sim 0.5$ because of $v/c \sim 0.1-0.3$

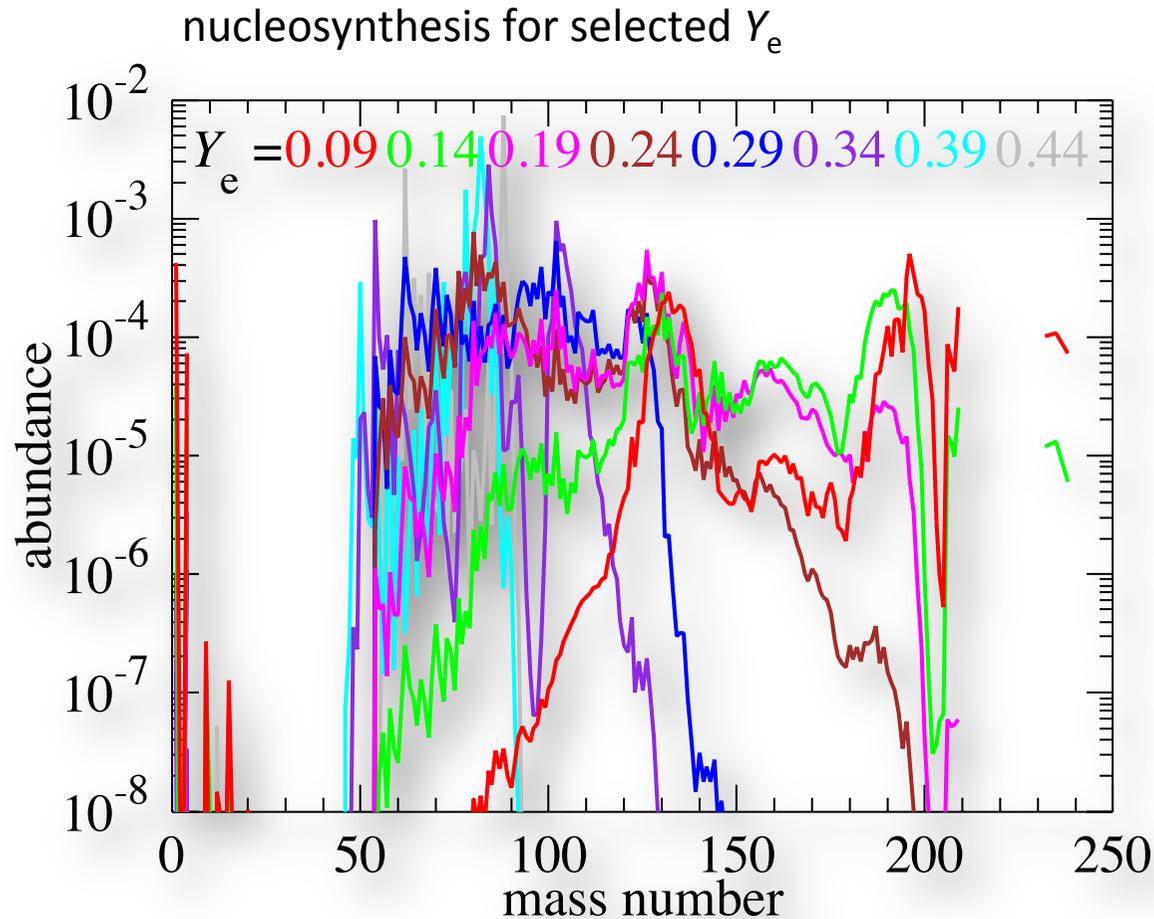


- ❖ higher and wider range of entropy per baryon ($= 0-50$) in contrast to previous cases ($= 0-3$)

$$Y_e = 0.09$$

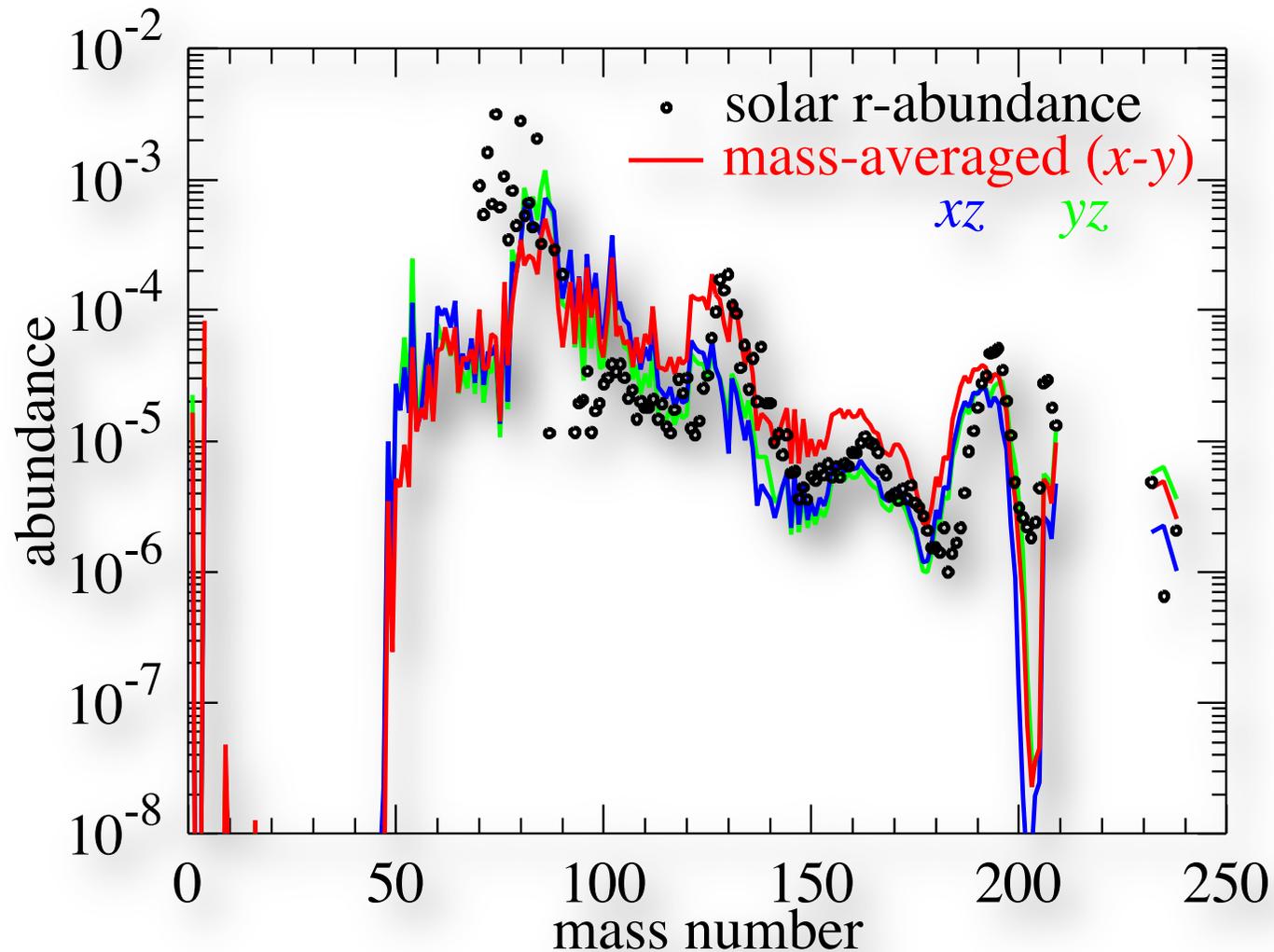


post-process nucleosynthesis



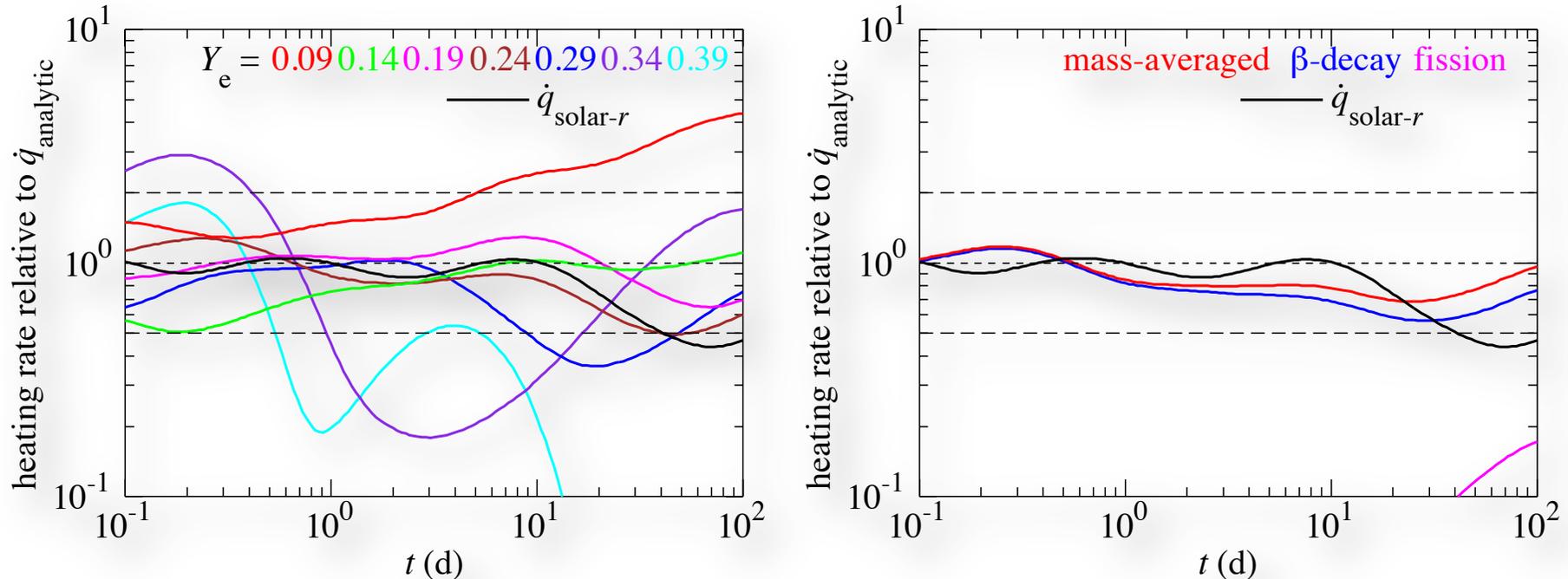
- ❖ variation of r-processes depending on Y_e
- ❖ production of iron to uranium
- ❖ no fission recycling

mass-integrated abundances



❖ reasonable agreement with full solar r-process range for $A = 90-240$

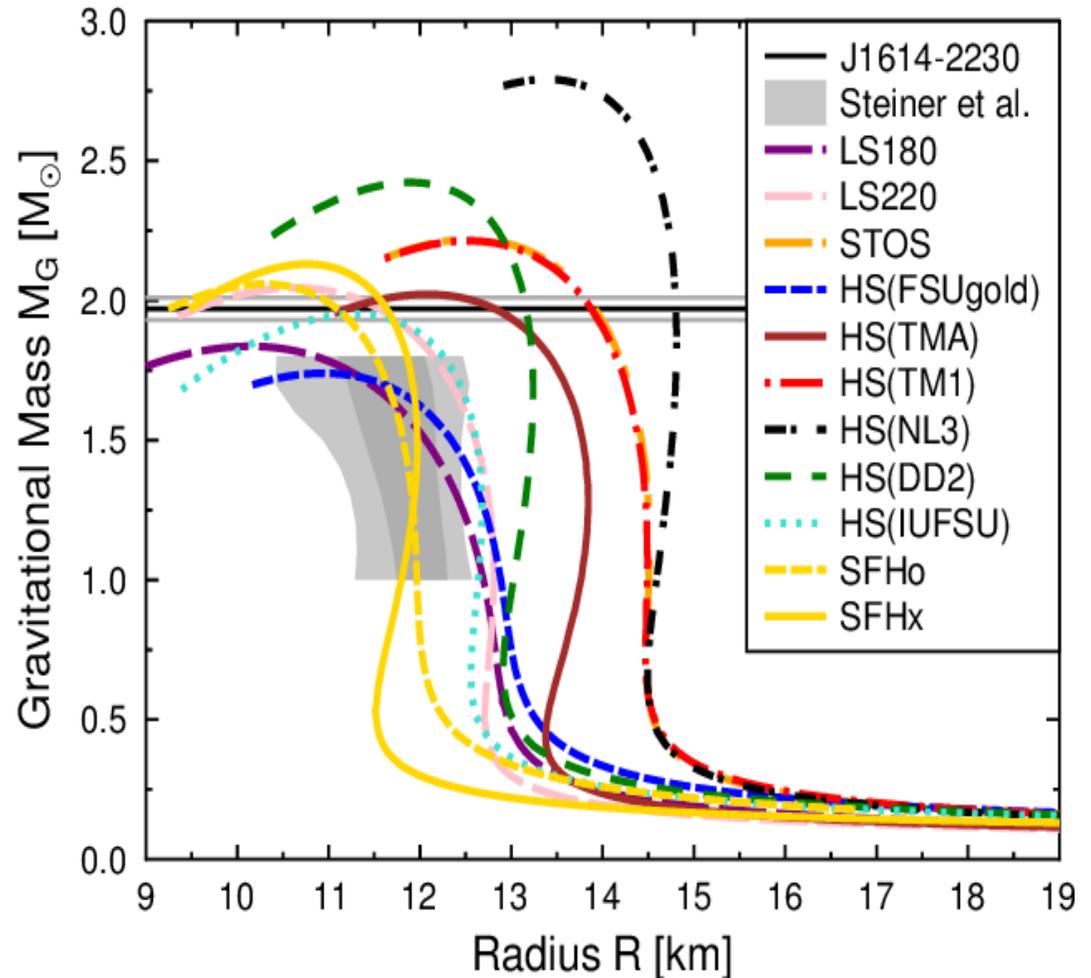
heating rate for the NS-NS ejecta



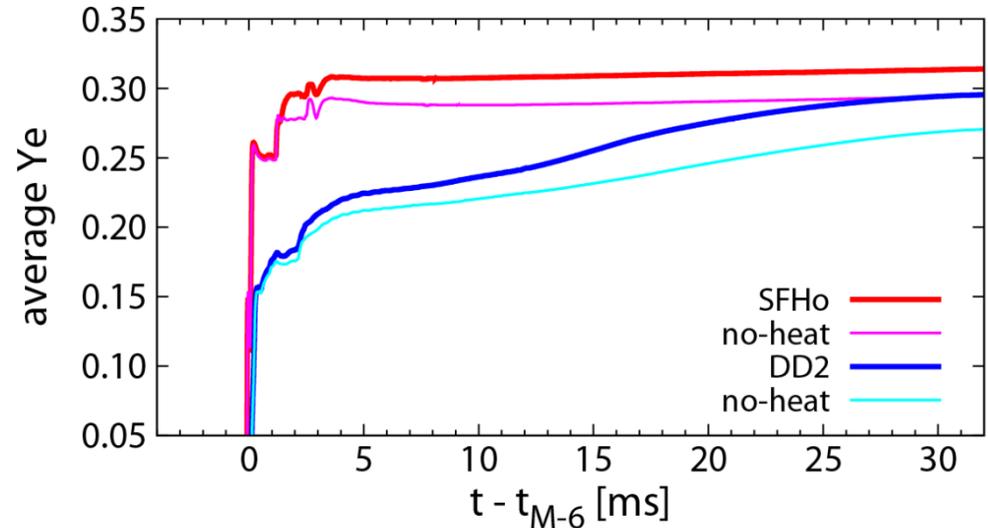
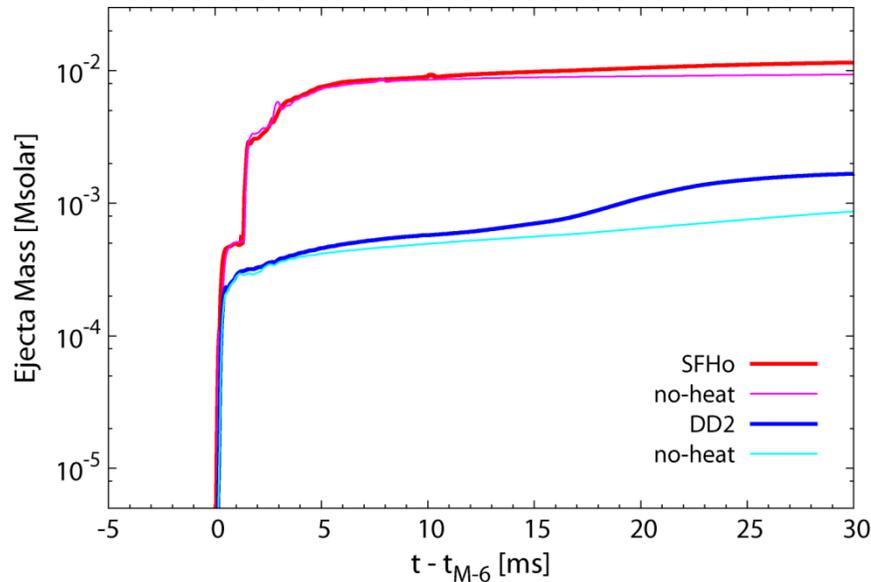
- ❖ heating rate for the mass-averaged abundances well fitted by the scaling law $dq/dt \sim t^{-1.3}$ (as well as by the solar r-pattern case)
- ❖ but dependent on Y_e ; there might be directional (polar to equatorial) differences

Dynamical mass ejection mechanism & EOS

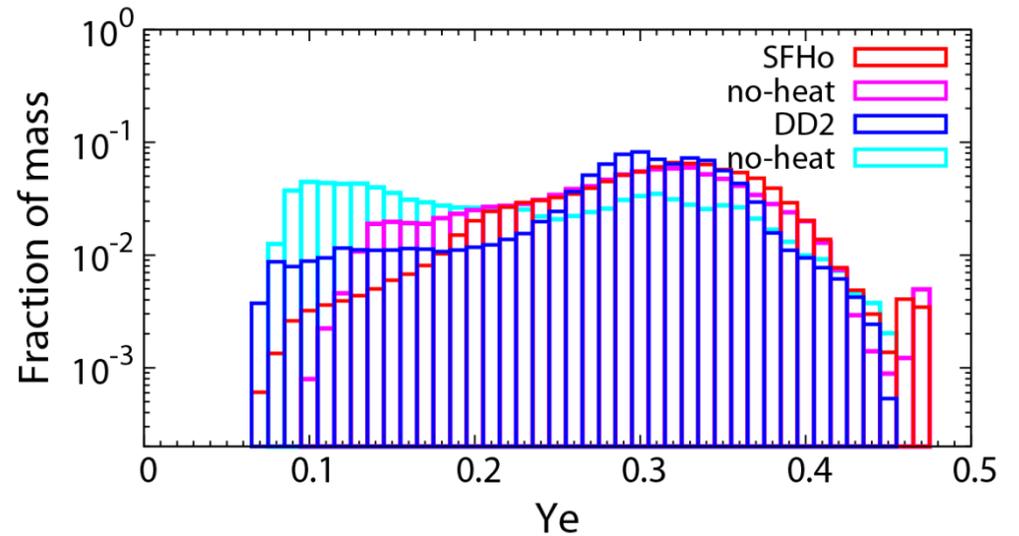
- ▶ ‘Stiffer EOS’
 - ▶ **TM1, TMA**
 - ▶ R_{NS} : lager
 - ▶ Tidal-driven dominant
 - ▶ **Ejecta consist of low T & Y_e NS matter**
- ▶ ‘Intermediate EOS’
 - ▶ **DD2**
- ▶ ‘Softer EOS’
 - ▶ **SFHo, IUFSU**
 - ▶ R_{NS} : smaller
 - ▶ Tidal-driven less dominant
 - ▶ Shock-driven dominant
 - ▶ **Ye can change via weak processes**



Effects of neutrino heating

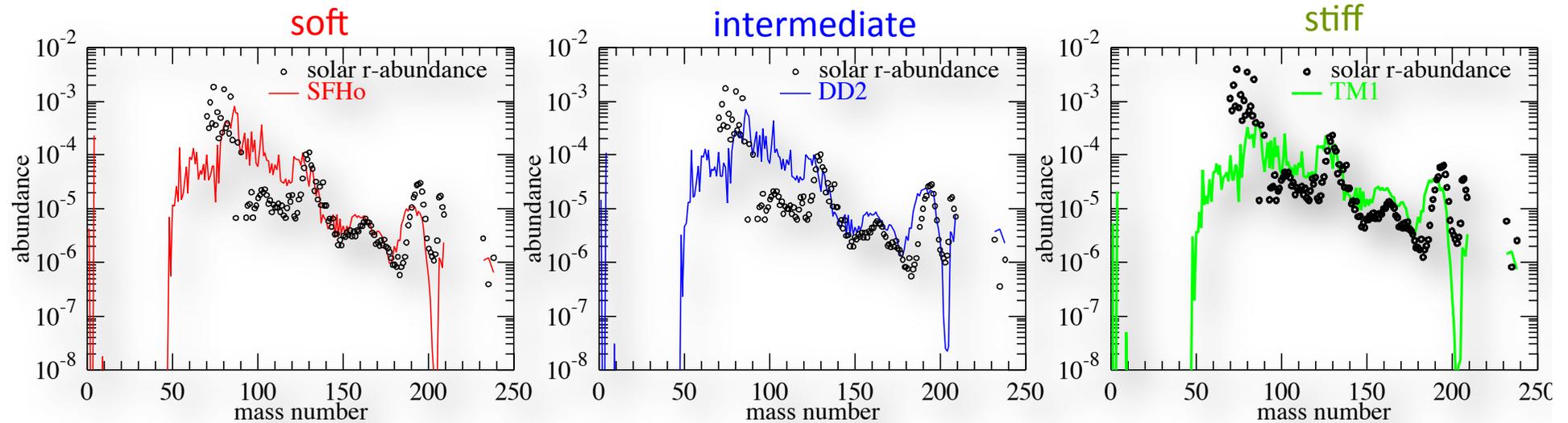


- ▶ Amount of ejecta mass can be increased $\sim 10^{-3}$ Msun
- ▶ Average Ye can change 0.02~0.03 depending on EOS : effect is stronger for stiffer EOS where HMNS survive in a longer time



dependence on EOSs

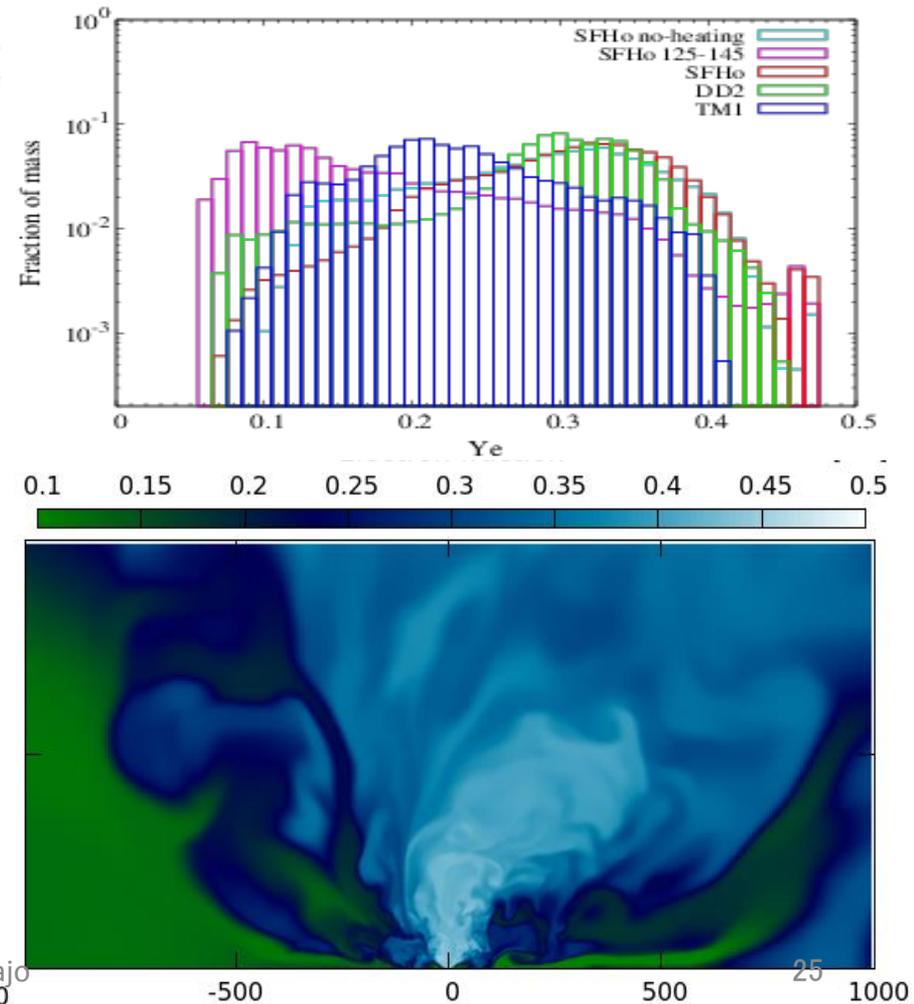
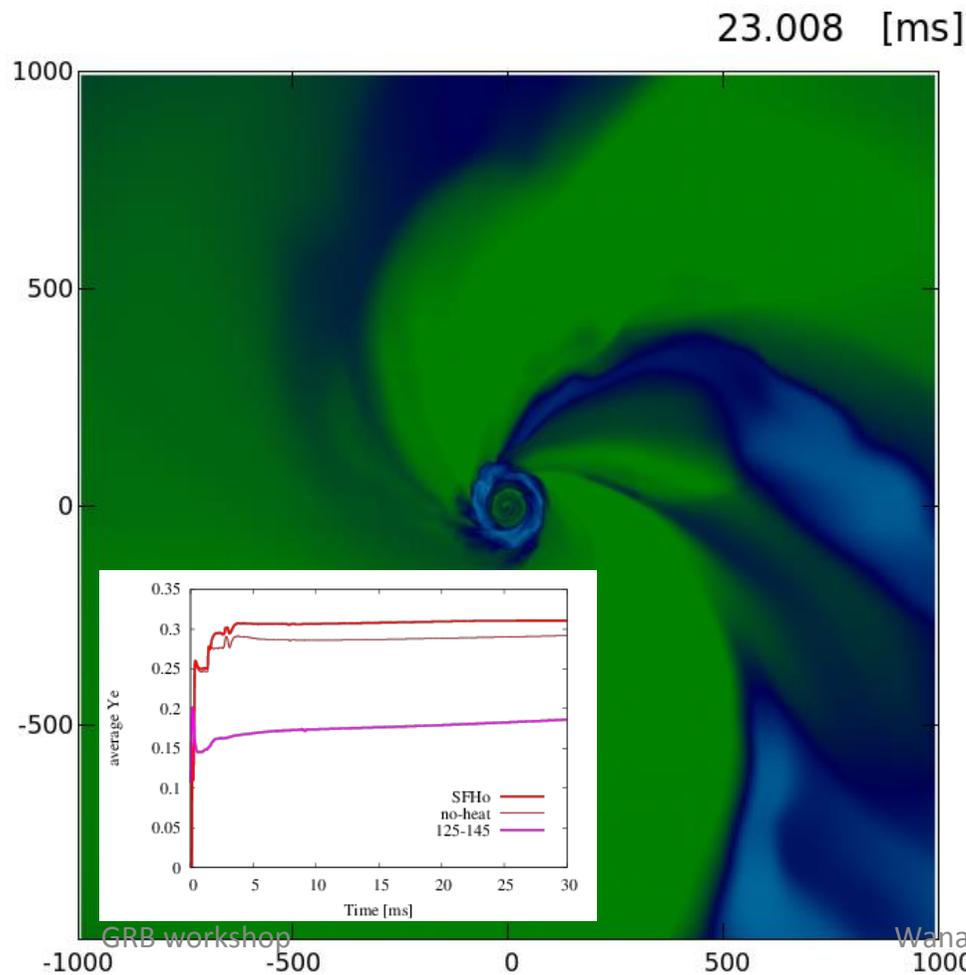
adopting nucleosynthesis of Wanajo+2014



- ❖ softer EOS predicts less heavy r-process products, but
- ❖ effects of EOSs are not large (good for the universality?)

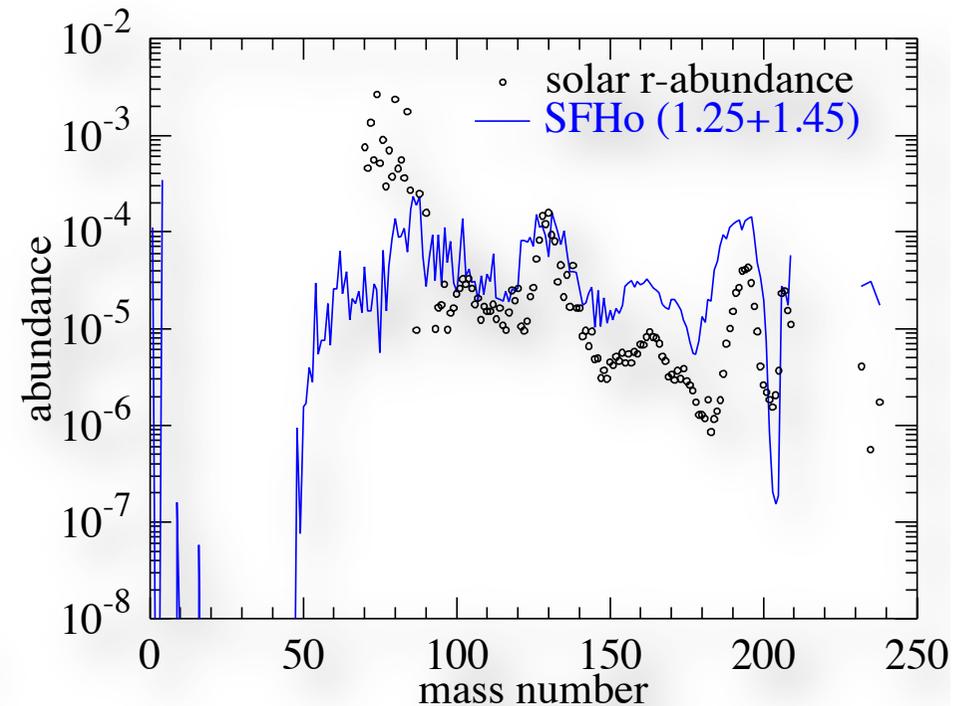
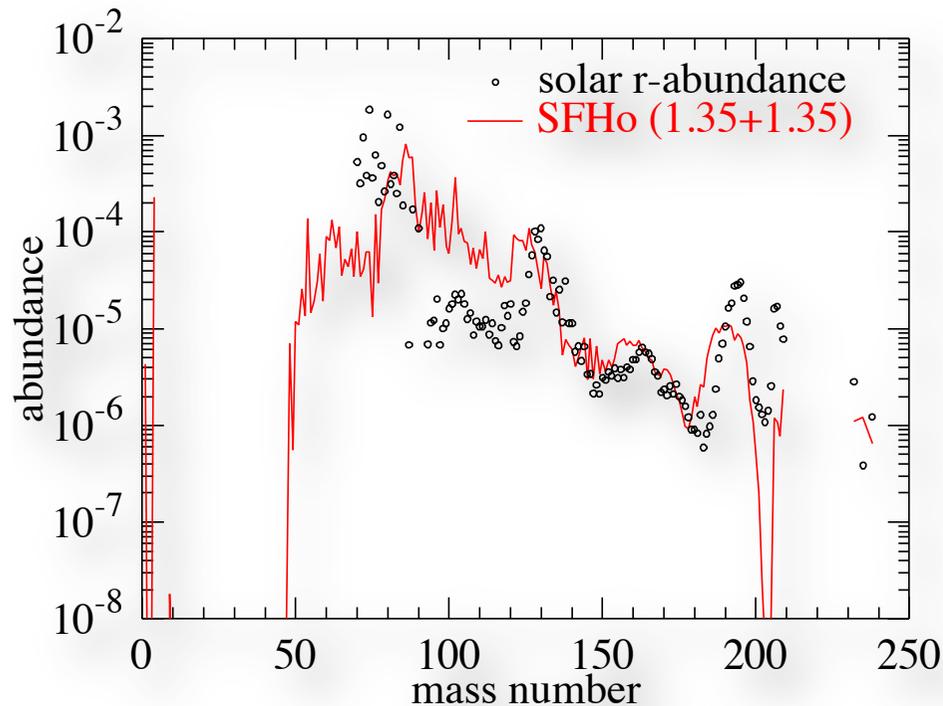
Unequal mass NS-NS system: SFHo1.25-1.45

- ▶ Orbital plane : Tidal effects play a role, ejecta is neutron rich
- ▶ Meridian plane : shock + neutrinos play roles, ejecta less neutron rich



dependence on the NS mass ratio

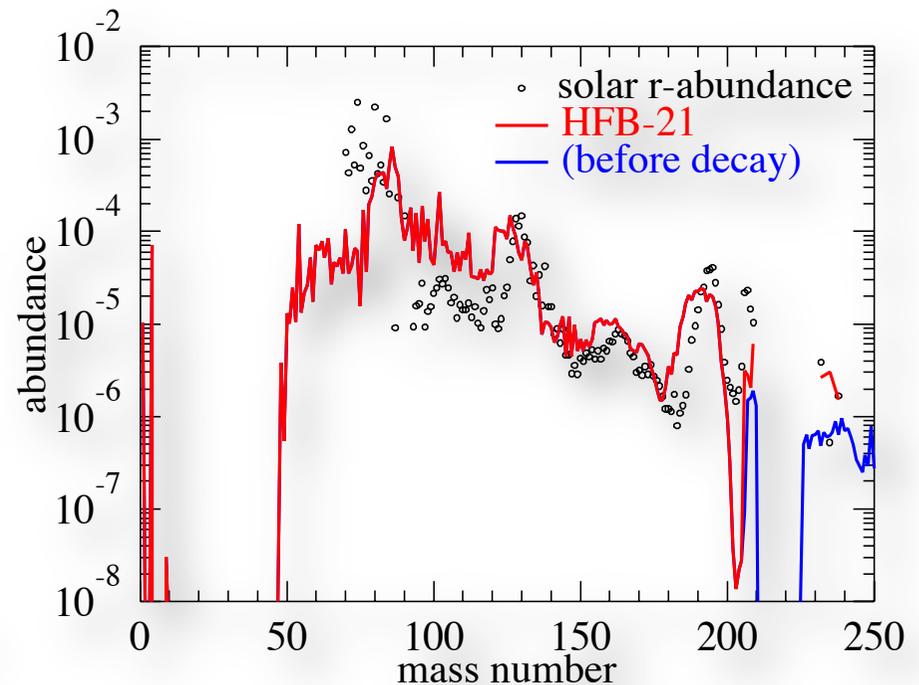
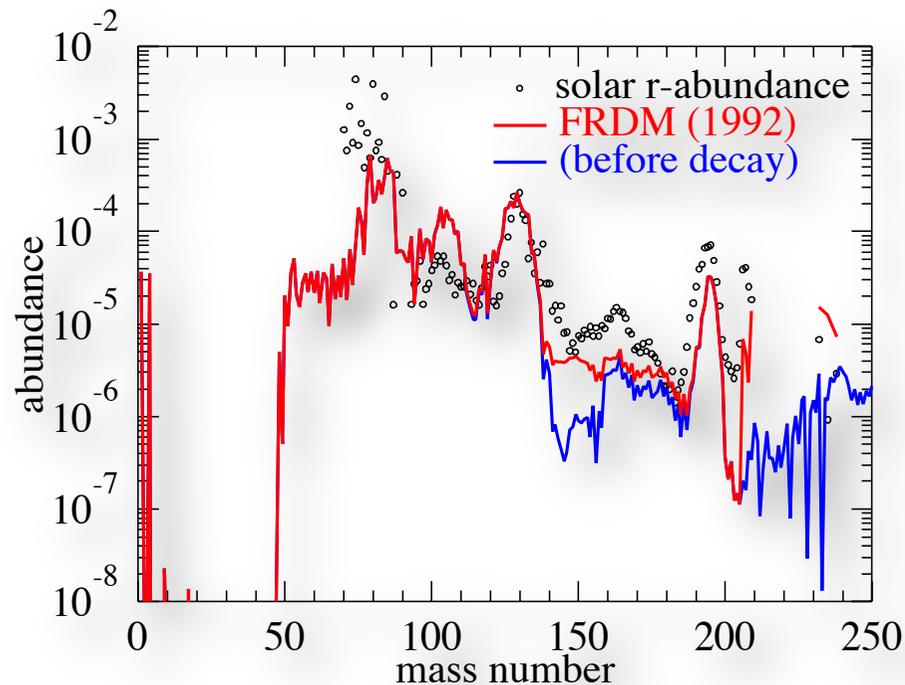
adopting nucleosynthesis of Wanajo+2014



- ❖ small asymmetry predicts less heavy r-process products
- ❖ moderate asymmetry is the best? (e.g., 1.3+1.4)

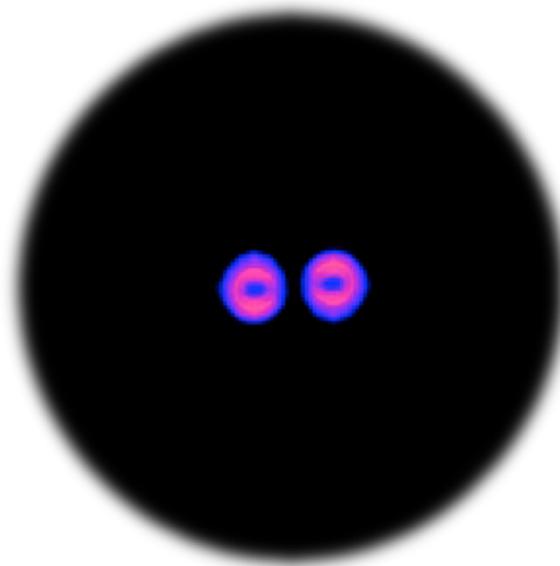
comparison for different mass models

for neutron star mergers in Wanajo+2014; without fission



❖ large differences between FRDM (1992, **not 2012!**) and HFB-21

summary and outlook



- ❖ NS mergers: very promising site of r-process and sRGBs
 - neutrinos play a crucial role (in particular for a soft EOS)
- ❖ still many things yet to be answered...
 - dependence on mass ratios of NSs and EOSs; how about BH-NS?
 - how the subsequent BH-tori contribute to the r-abundances?
 - r-process nova light curves as functions of time and viewing angle?