

# NS physics with GW observations: the importance of being magnetic

**Simone Dall'Osso**

**TAT - University of Tübingen**

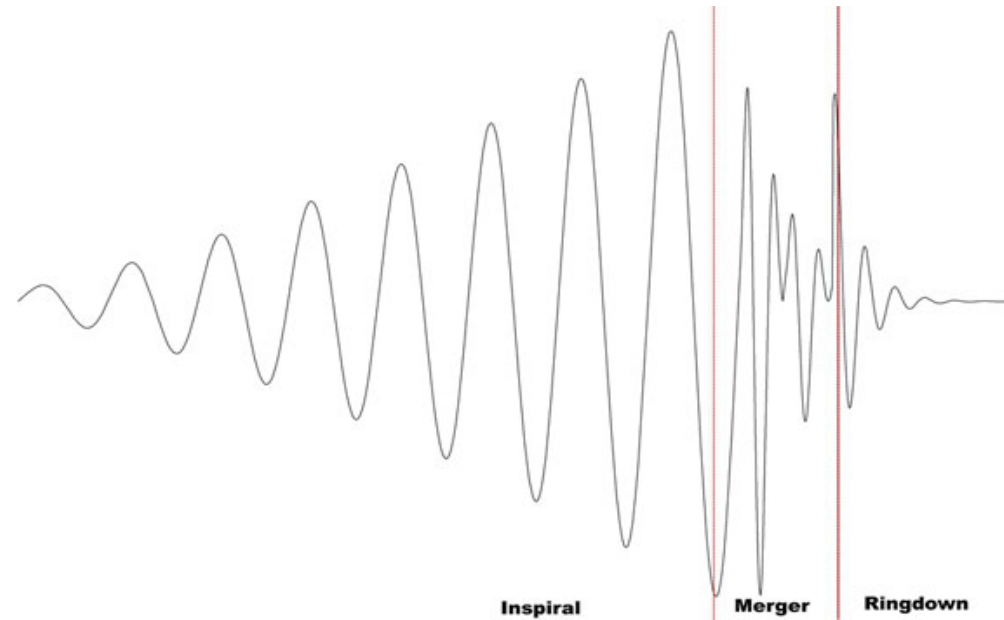
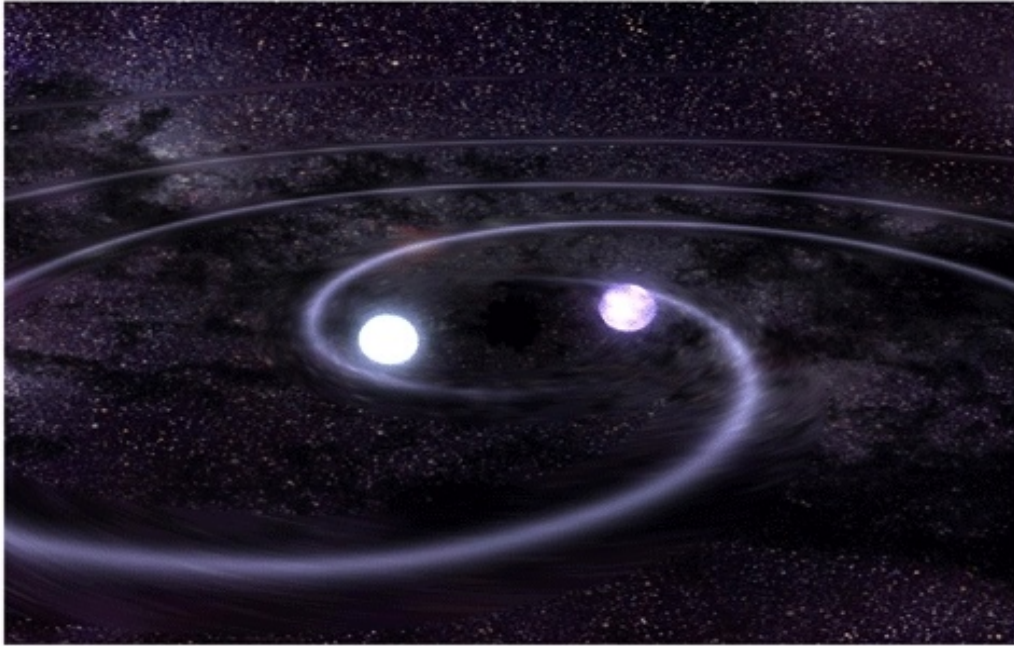
in collaboration with

Bruno Giacomazzo (Trento University)

Rosalba Perna (Stony Brook University)

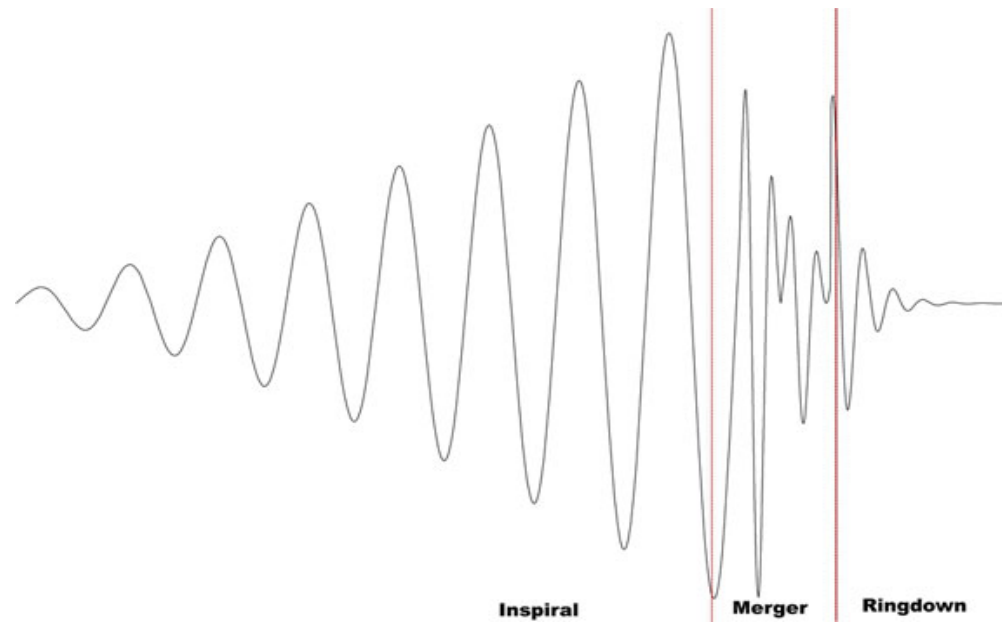
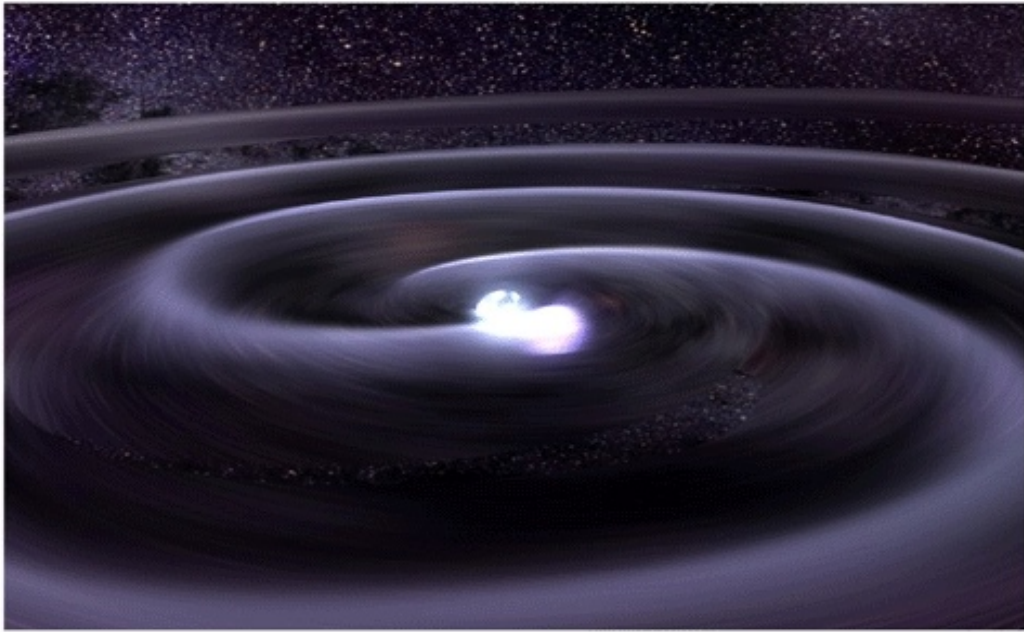
Luigi Stella (Astronomical Observatory of Rome)

# Binary NS mergers



inspiral  $\longrightarrow$  GW-driven,  $\sim$  point-mass approx

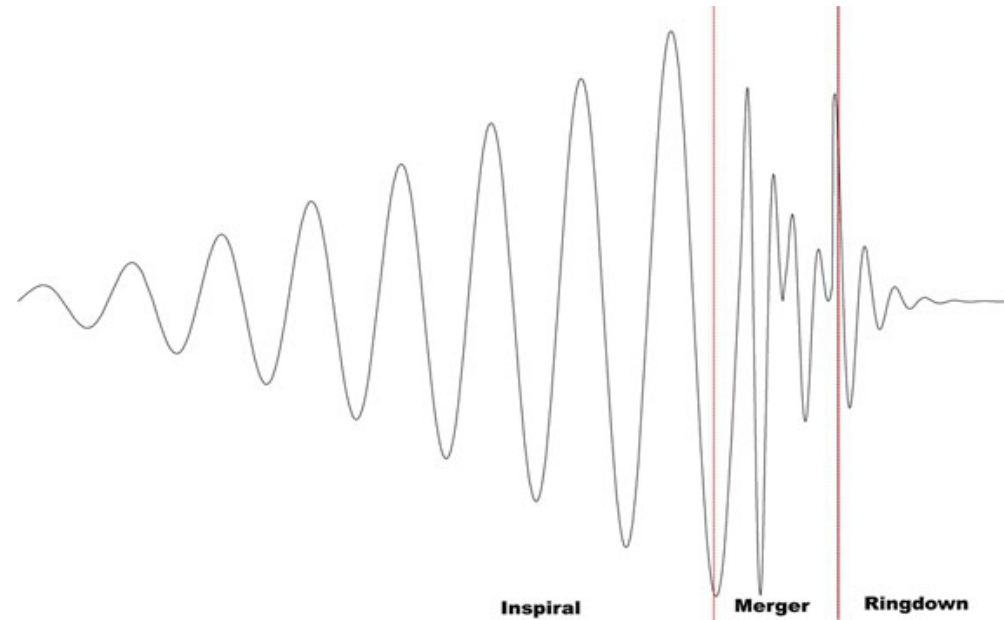
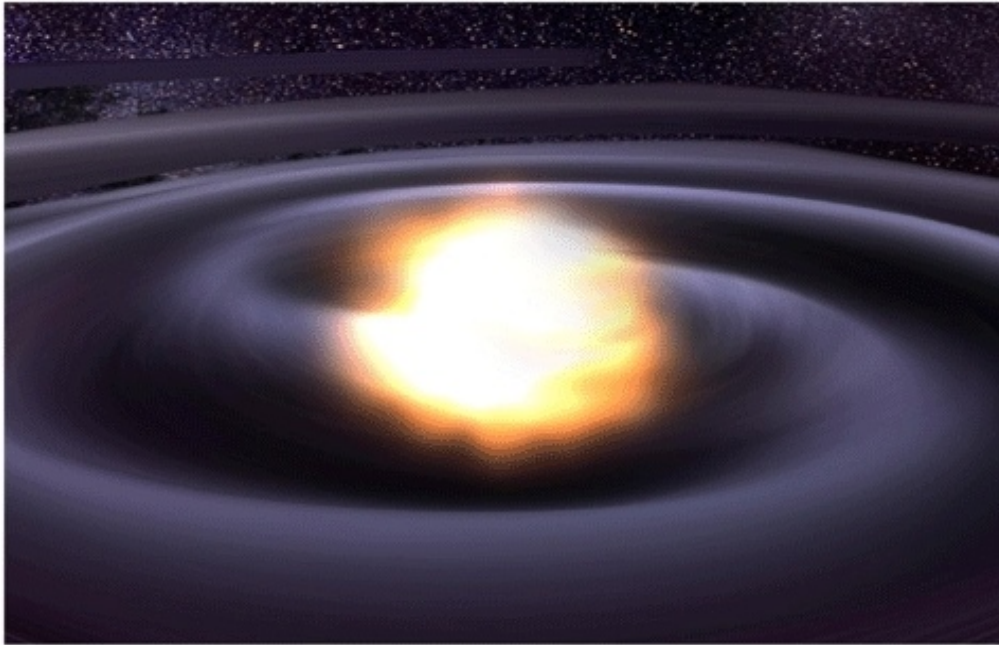
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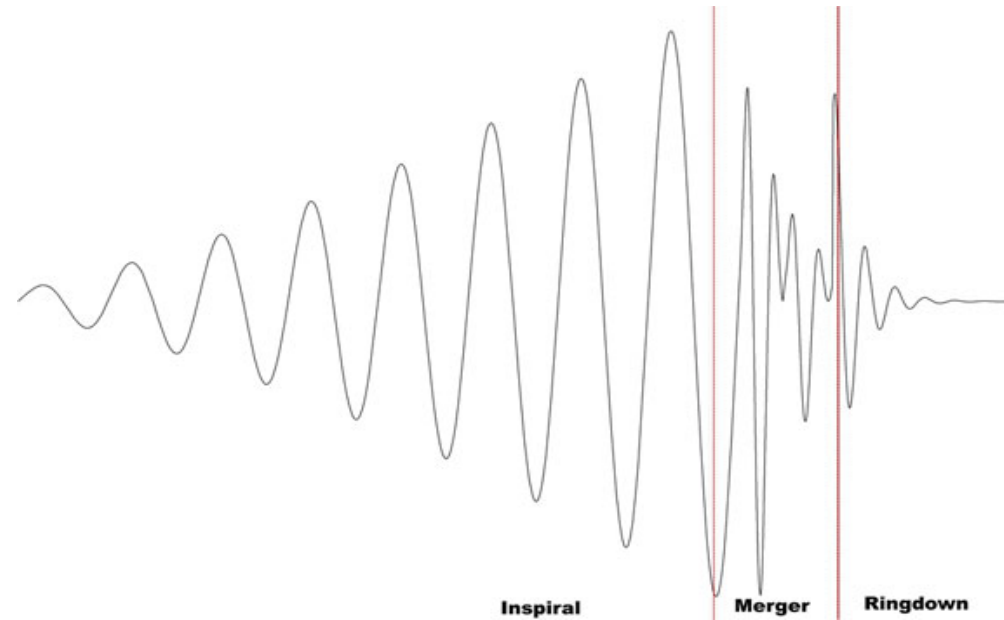
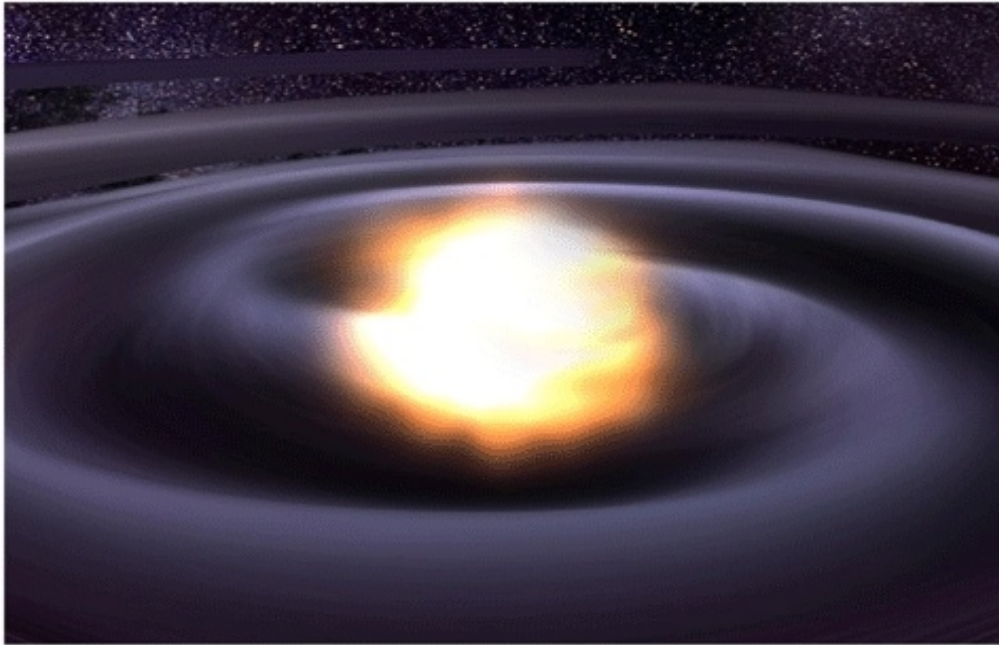


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 ↖ Short GRBs   
 ↘ GW signals

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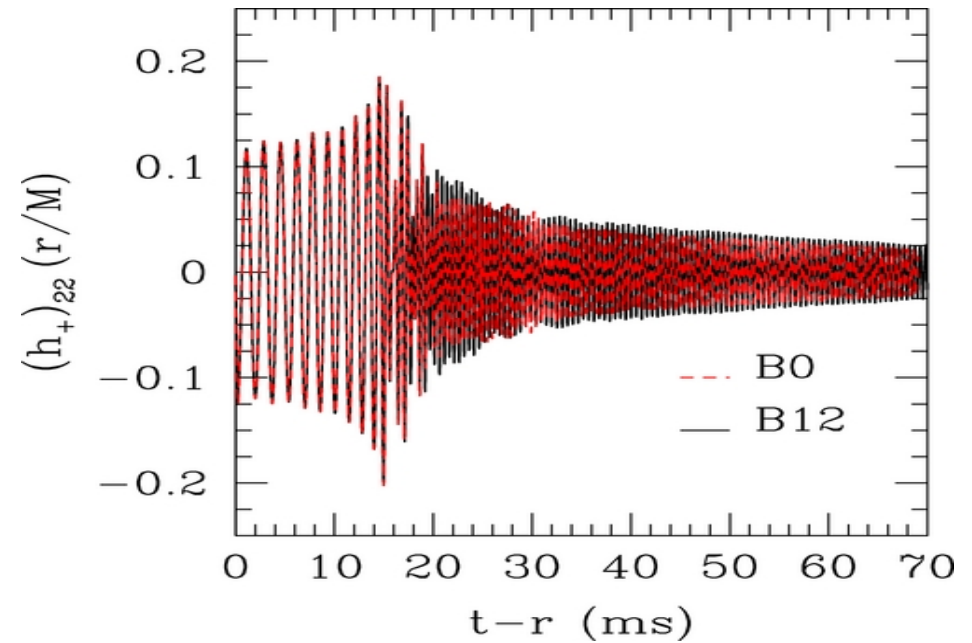
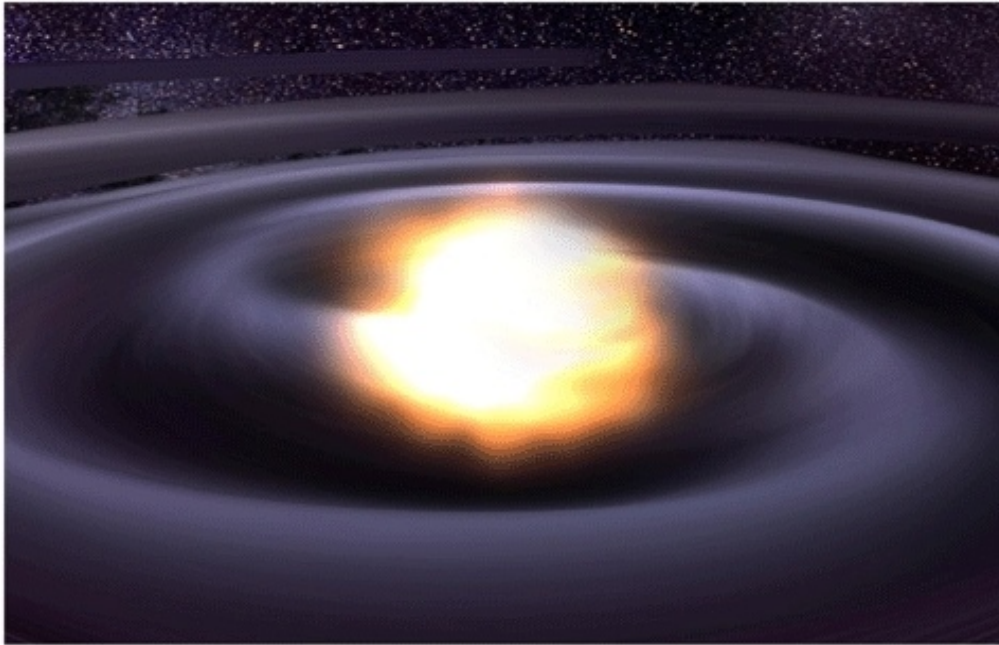


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Giacomazzo & Perna (2013)

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# Post-merger remnants

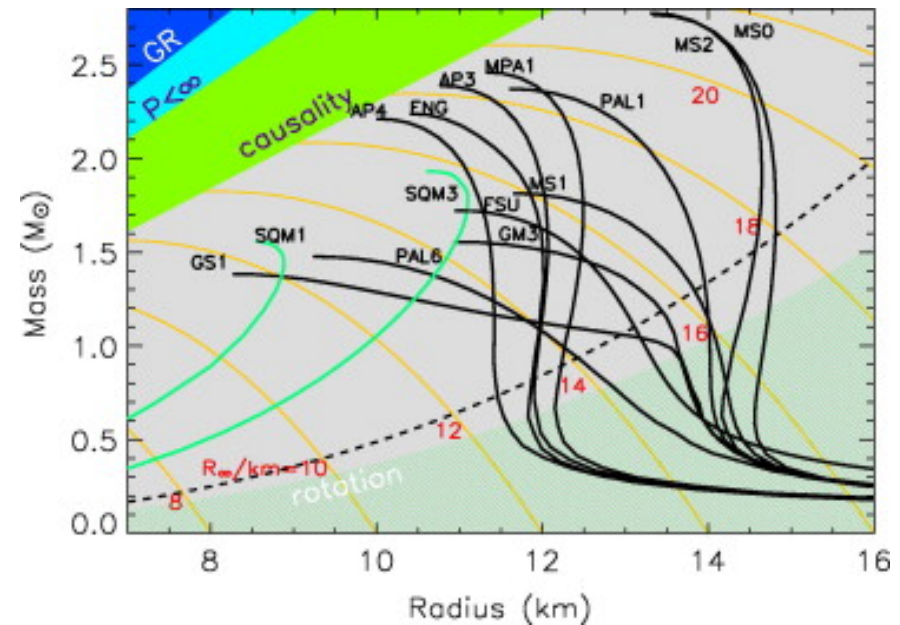
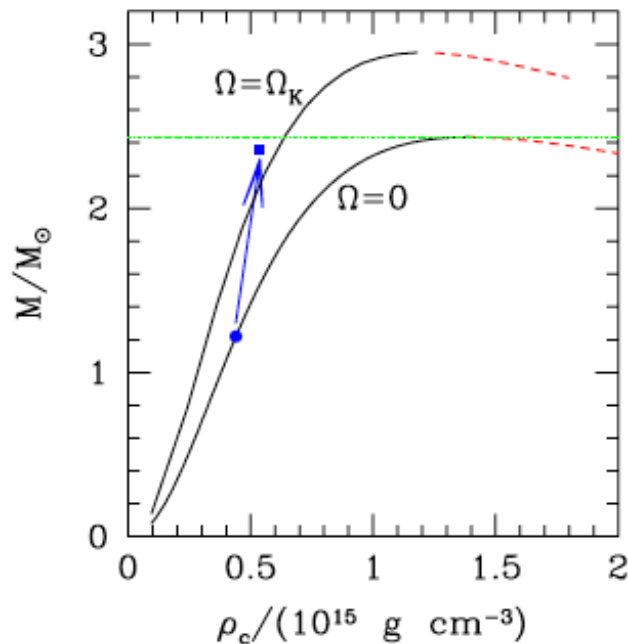
- Direct collapse to BH if  $M_{\text{TOT}} > M_{\text{max}}(\Omega)$
- Formation of an unstable NS if  $M_{\text{max}}(\Omega) > M_{\text{TOT}} > M_{\text{max}}$
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$$M_{\text{max}}(\Omega) = M_{\text{max}}(1 + \alpha \Omega^\beta)$$

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# Post-merger remnants

Differential rotation/turbulence → strongly twisted internal field

$$E_B \geq 10^{50} \text{ erg}$$

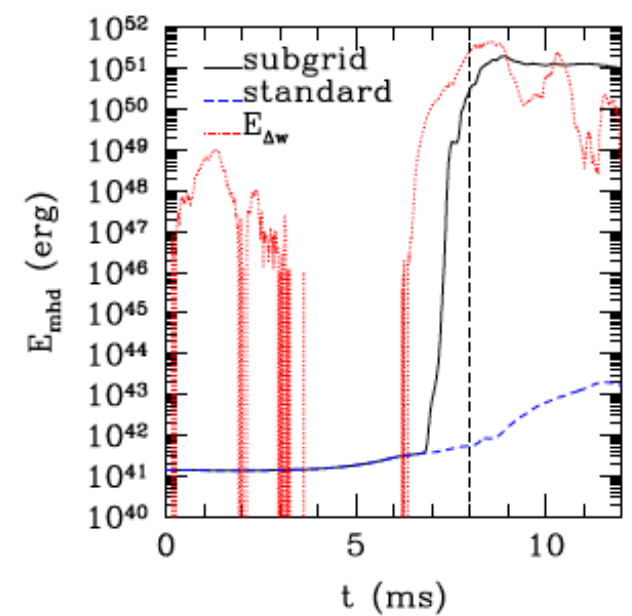
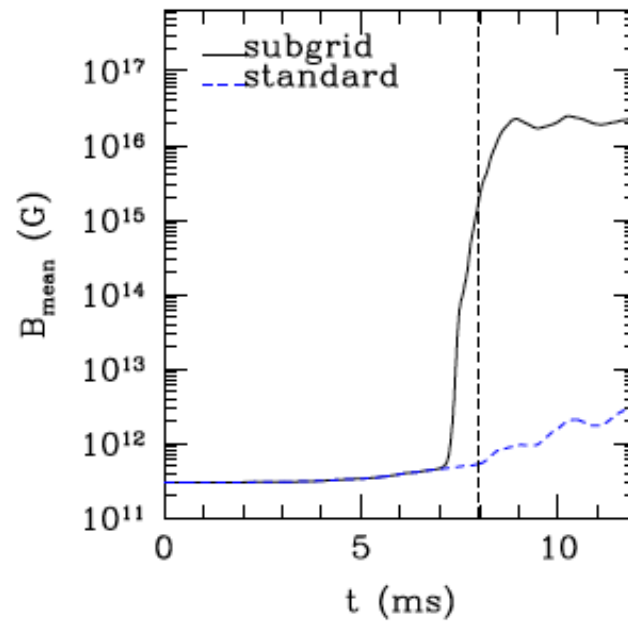
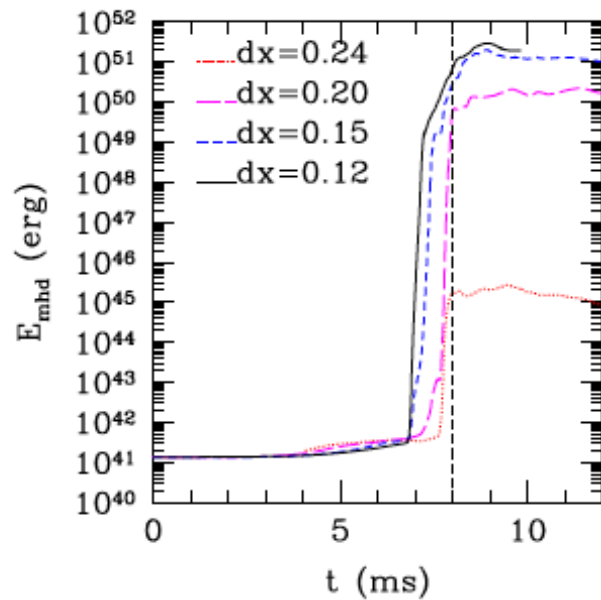
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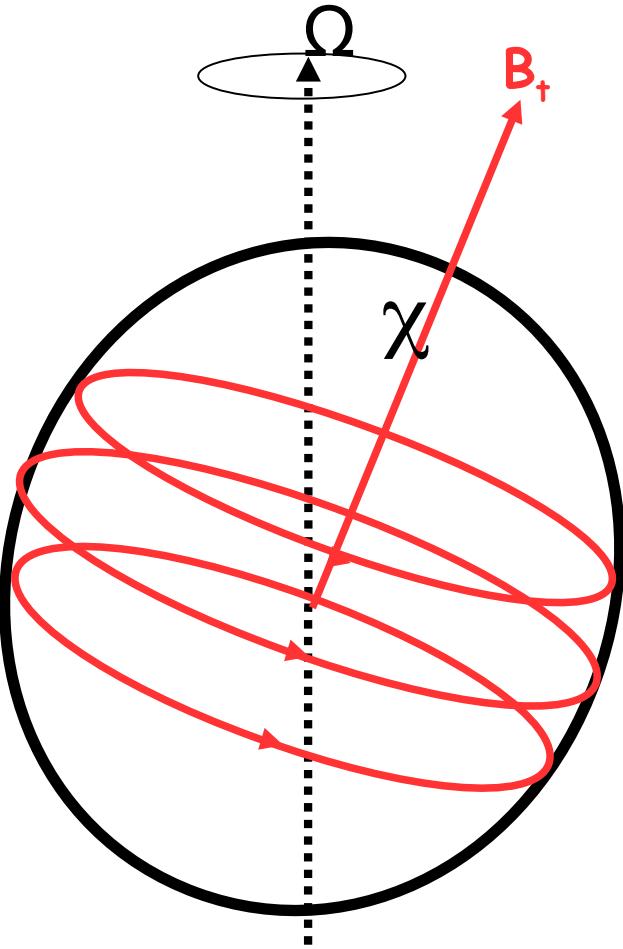
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# GW emission: a generic mechanism

TOROIDAL B  $\square$

PROLATE DISTORTION



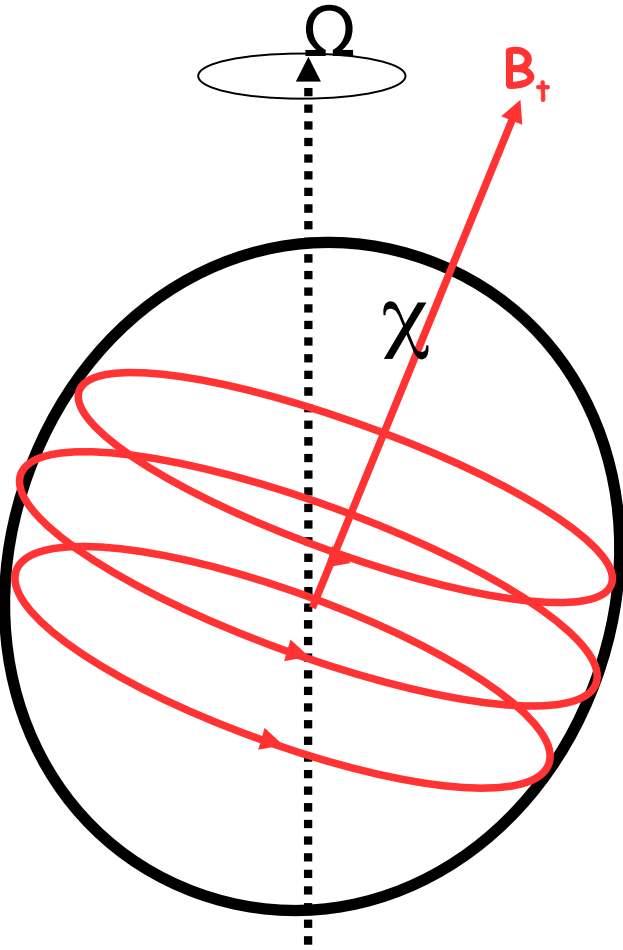
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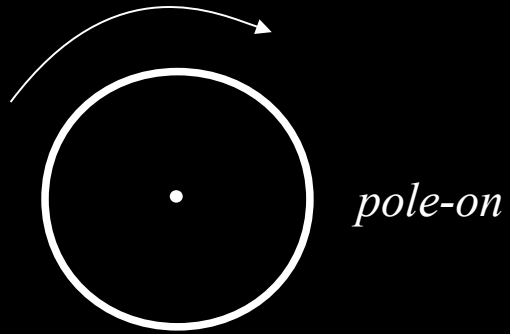
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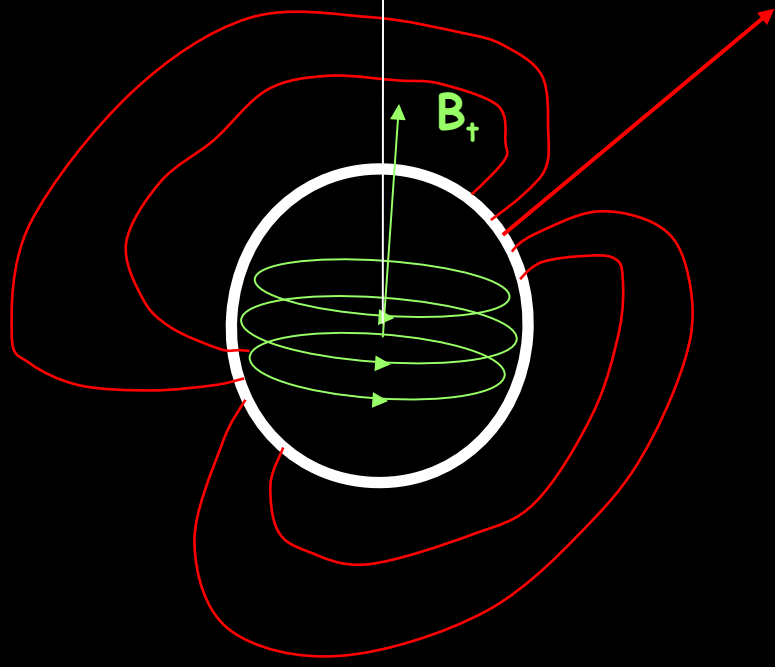
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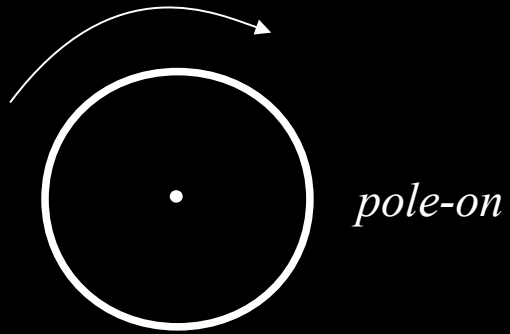
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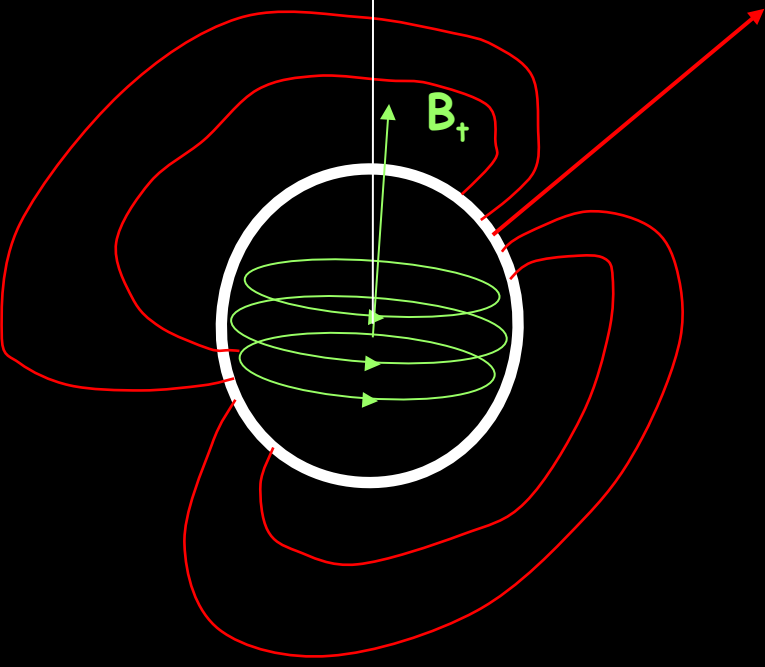


**spin**





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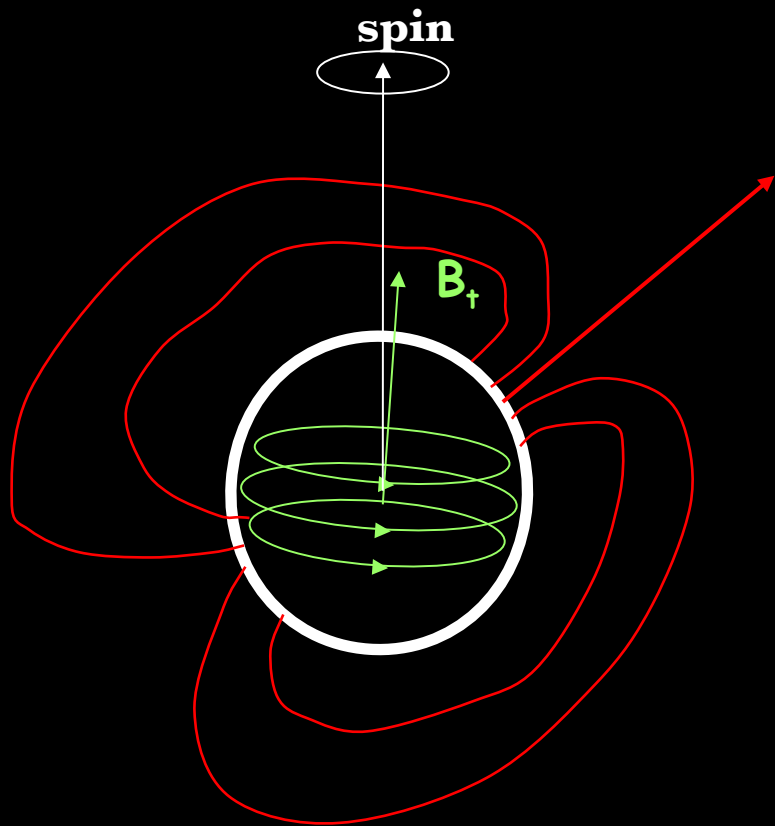
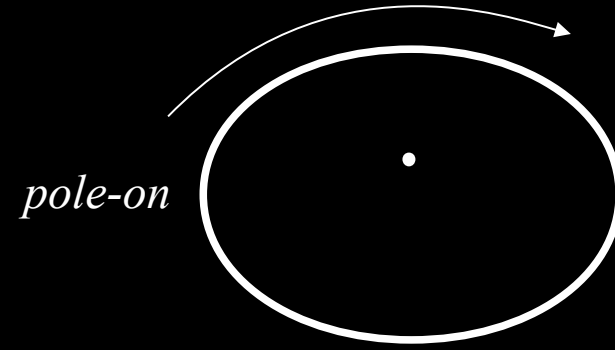
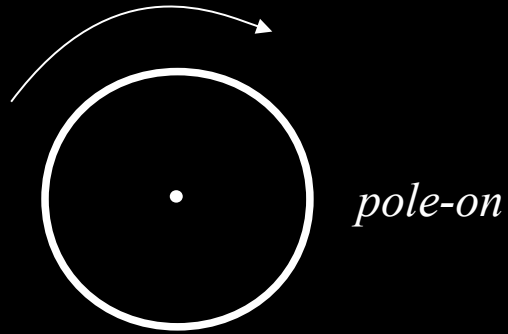


**Internal  
dissipation**



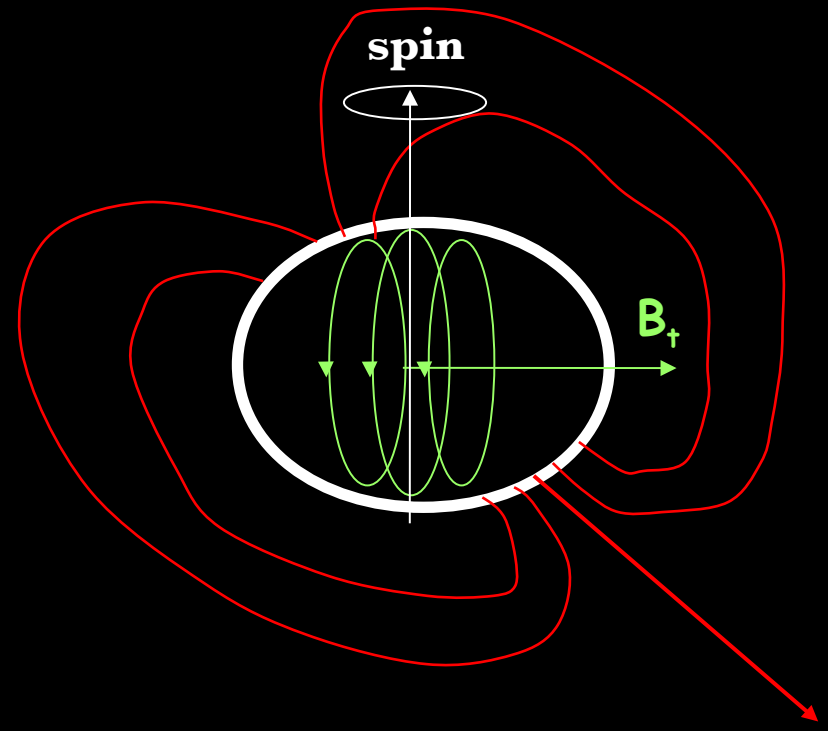
$$\mathbf{L} = \mathbf{I}\boldsymbol{\Omega}$$
$$\mathbf{E} = \mathbf{I}\boldsymbol{\Omega}^2$$

*GW emission  
maximised!*



Internal  
dissipation  
→

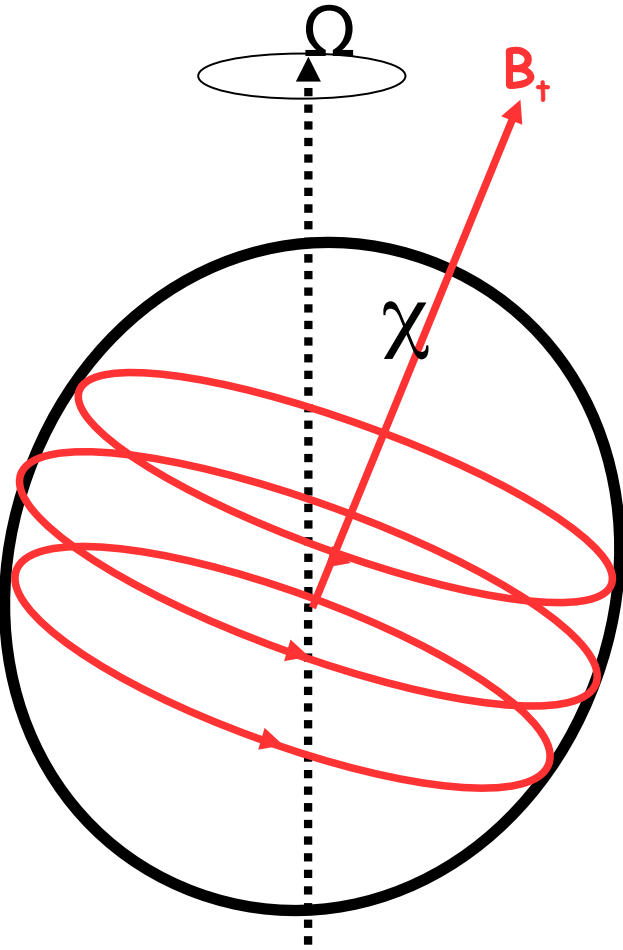
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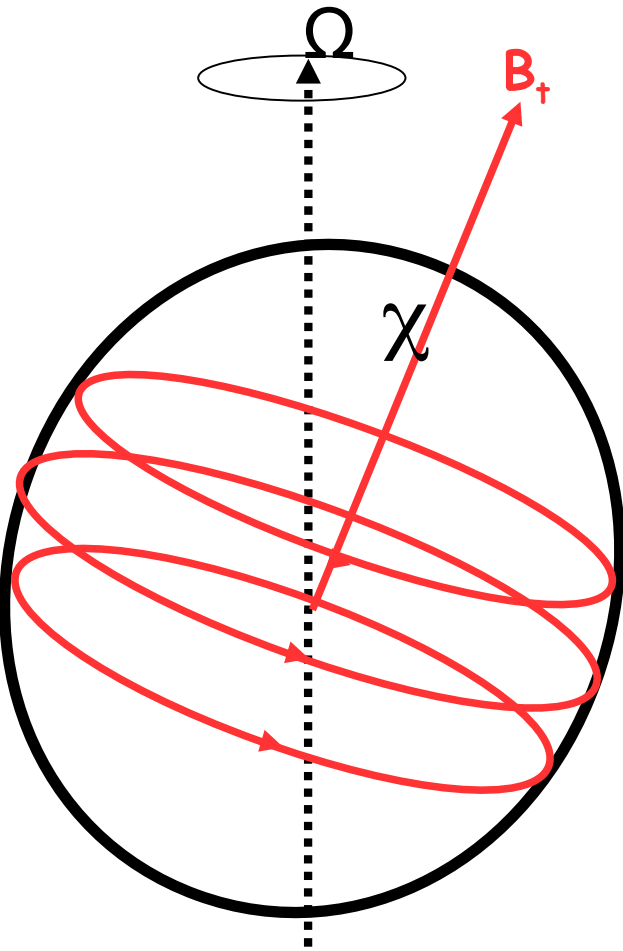
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Proposed application in different contexts  
where differential rotation plays a role

## **1. magnetar formation in core-collapse**

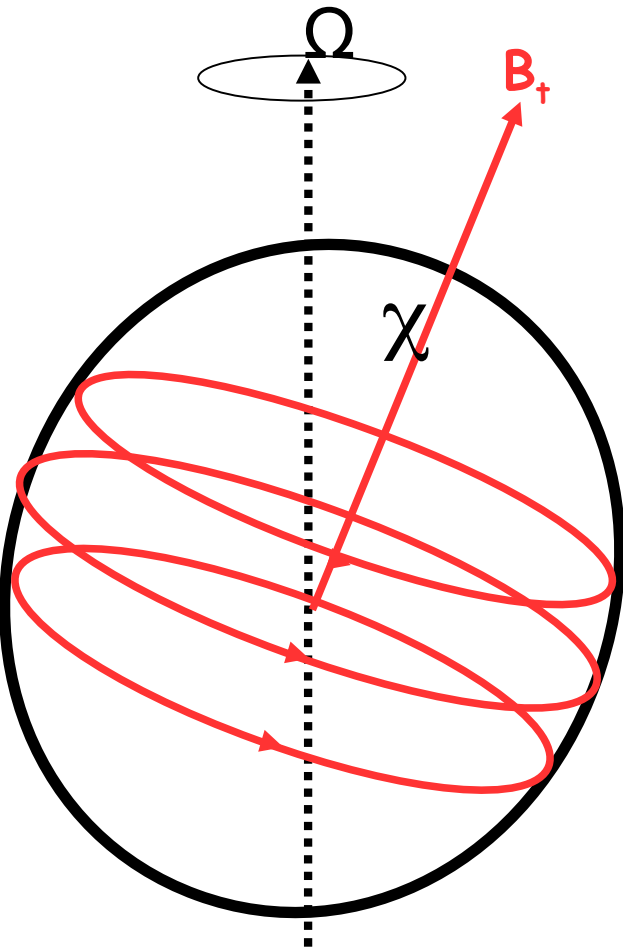
Thompson & Duncan 1992, 1993

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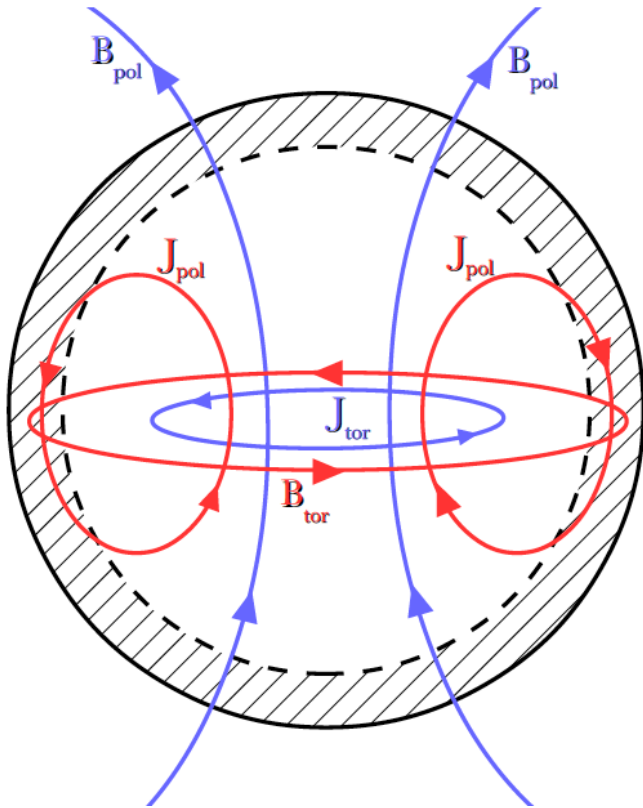
Stella et al. 2005, Dall'Osso, Shore & Stella 2009

## 2. *r*-modes in accreting NSs in LMXBs

Rezzolla et al. (2001)

Cuofano et al. (2012)

# Interior vs. exterior field: stability

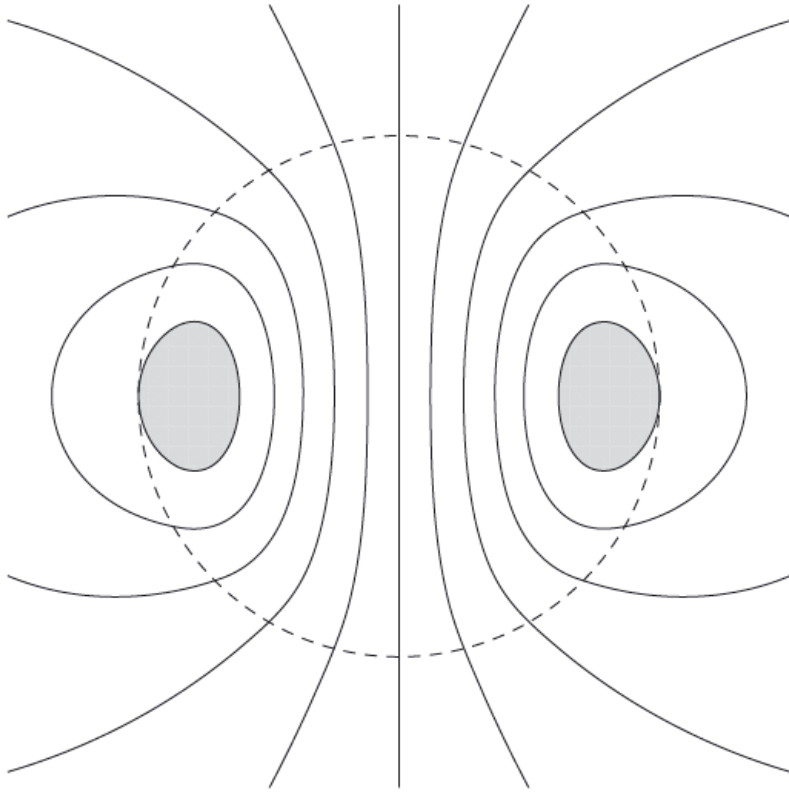


So-called twisted-torus is a generic relaxed state for NS global fields (Spruit & Braithwaite 2006)

Stability limit for such configurations depends critically on stable stratification of NS matter (Reisenegger 2009, Akgün et al. 2013)

Large toroidal-to-poloidal field ratios can be stable (Braithwaite 2009, Akgün et al. 2013)

# Magnetically-induced ellipticity



$$\epsilon_B^0 \sim \frac{E_B}{E_{grav}} \simeq 3.5 \times 10^{-4} B_{16}^2 R_6^4 M_{1.4}^{-2}$$

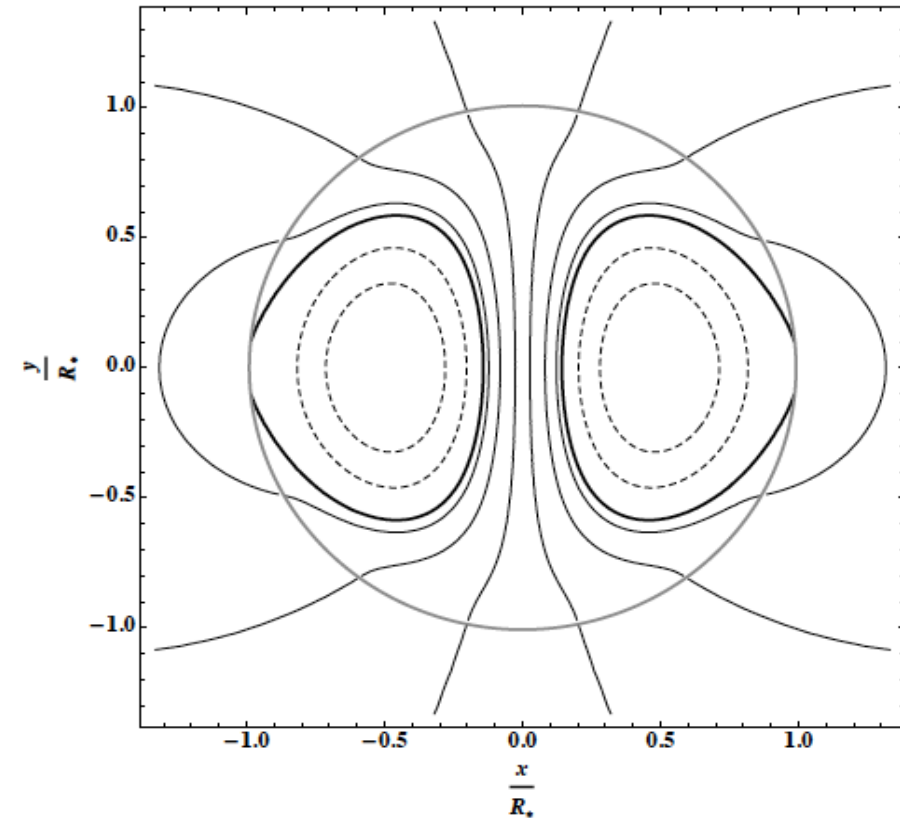
$$\epsilon_B \simeq 3.4 \times 10^{-7} B_{dip,14}^2 \frac{R_{12}^4}{M_{1.4}^2} \left(1 - 0.37 \frac{E_T}{E_{pol}}\right)$$

$$\frac{E_T}{E_{pol}} < \frac{80}{B_{pol,14}} \frac{M_{2.36}}{R_{12km}^2}$$

Mastrano et al. 2011

Akgün et al. 2013

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$$\epsilon_B^{new} \simeq 3.2 \times 10^{-6} B_{dip,14}^2 \frac{R_{12}^4}{M_{1.4}^2} \left(1 - 0.73 \frac{E_T}{E_{pol}}\right)$$

$$\frac{E_T}{E_{pol}} < \frac{500}{B_{pol,14}} \frac{M_{2.36}}{R_{12km}^2}$$

# Damping Time

Orthogonalization of the newly formed magnetar  
caused by internal dissipation of free precession

Must be faster than *em* spindown  
 $\tau_{\text{em}} \diamond 1.1 (P_{\text{ms}}/B_{\text{d},14})^2 \text{ days} \sim 10^5 \text{ s}$

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$$\tau_{\text{diss}} \sim 5 (B_t/10^{16}\text{G})^2 (P/\text{ms})^2 (T/10^{10}\text{K})^{-6} \text{ s}$$

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**Short dissipation timescale  $\sim 10^{2-3} \text{ s}$**



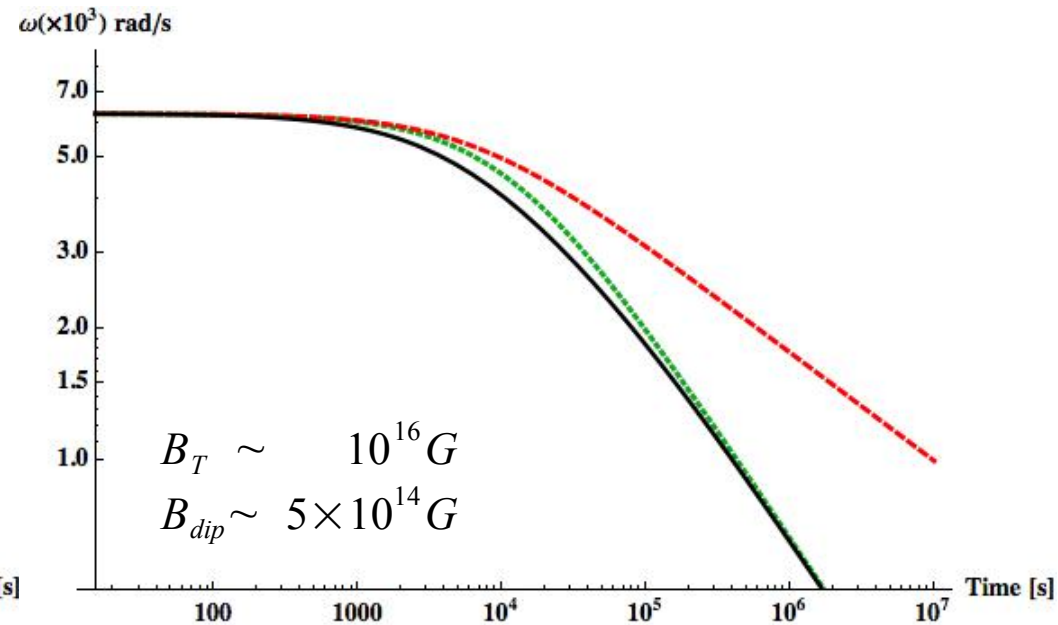
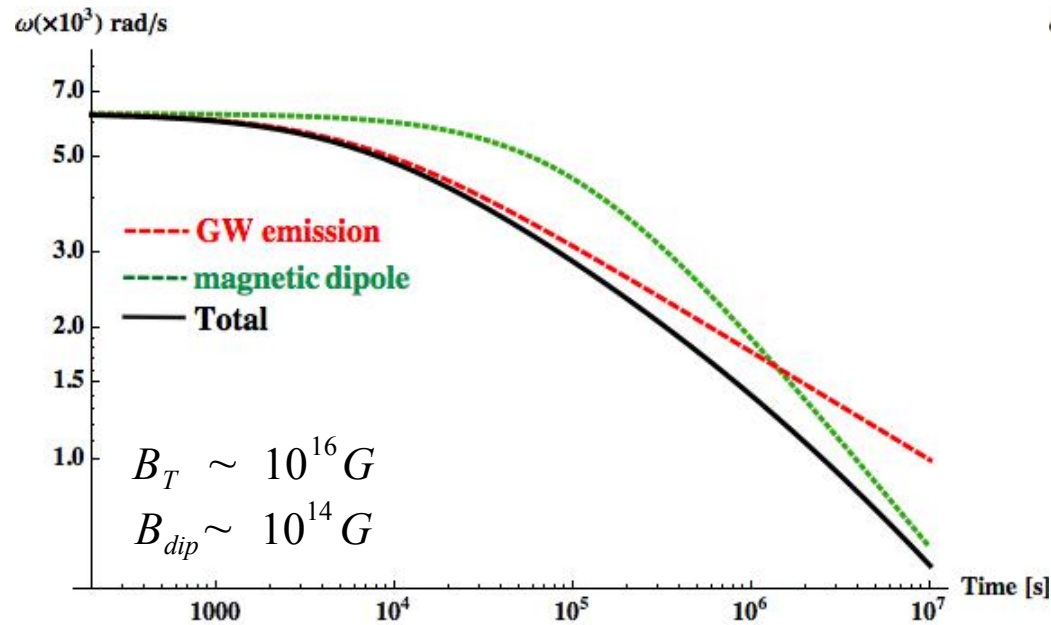
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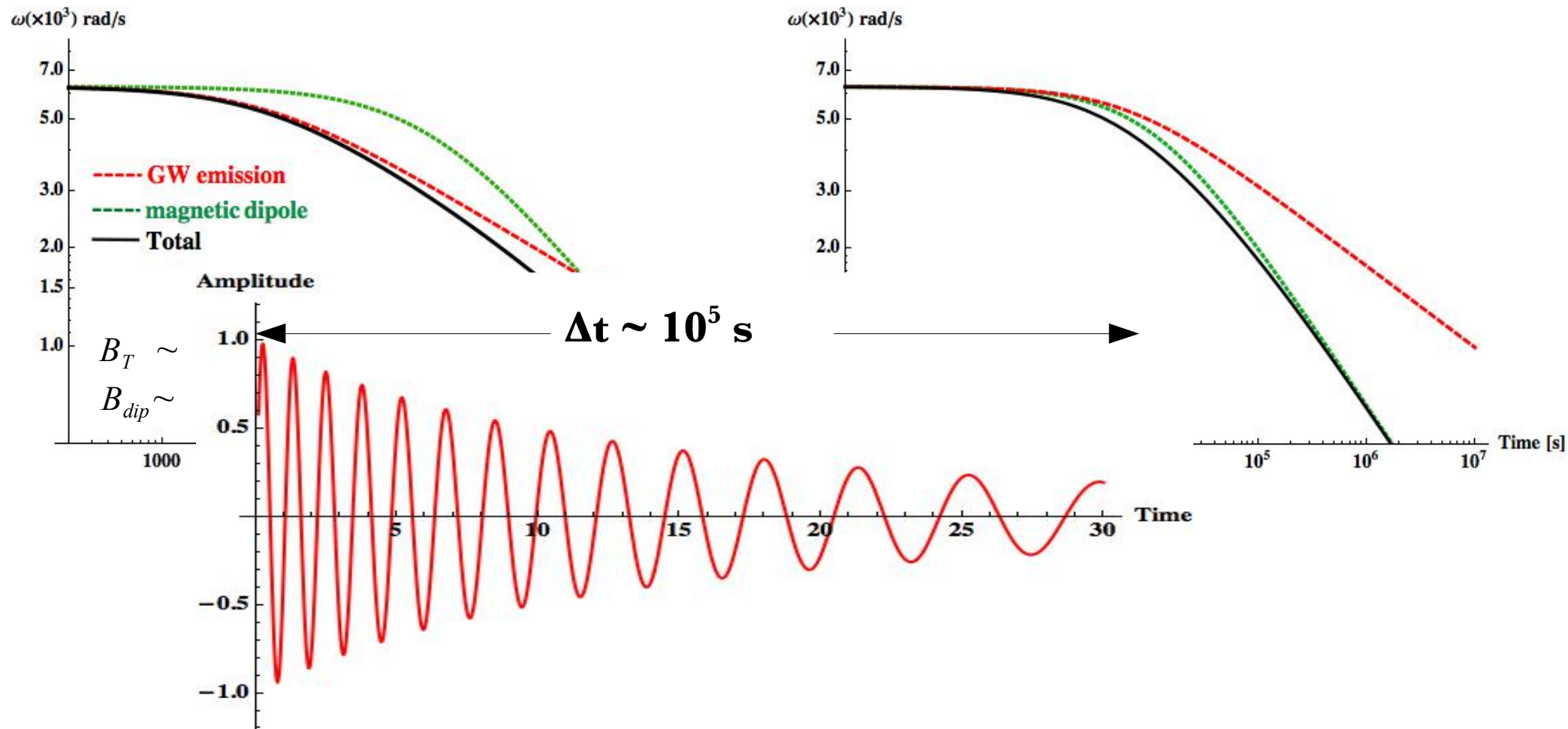
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$$S/N = 2 \left[ \int \frac{|\tilde{h}(f)|^2}{S_h(f)} df \right]^{1/2} \quad \mathbf{D}_{\text{hor}} \sim \mathbf{35 \text{ Mpc}}$$

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**$D_{\text{hor}} \sim 200 \text{ Mpc}$**

for NS-NS inspirals

(Abadie et al. 2010)

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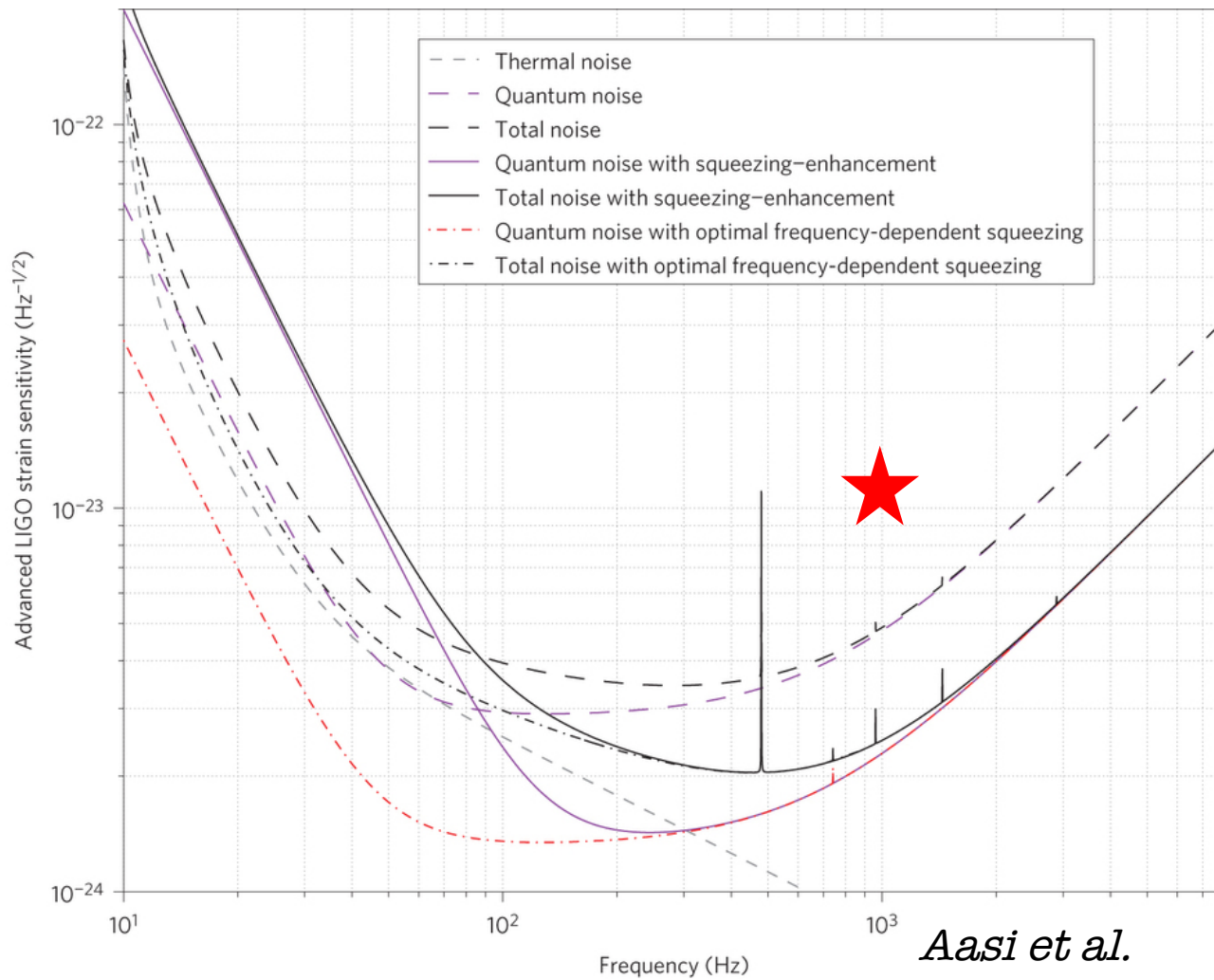
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$$\left( \frac{S}{N} \right)_{\text{GW}} \simeq 10 \left( \frac{D}{33.5 \text{Mpc}} \right)^{-1} \left( \frac{R}{15 \text{km}} \right) \left( \frac{M}{2.36 M_{\odot}} \right)^{1/2} \\ \times \left[ \left( \frac{f_f}{\text{kHz}} \right)^{-2} - \left( \frac{f_i}{\text{kHz}} \right)^{-2} \right]^{1/2} \quad \text{@35 Mpc}$$

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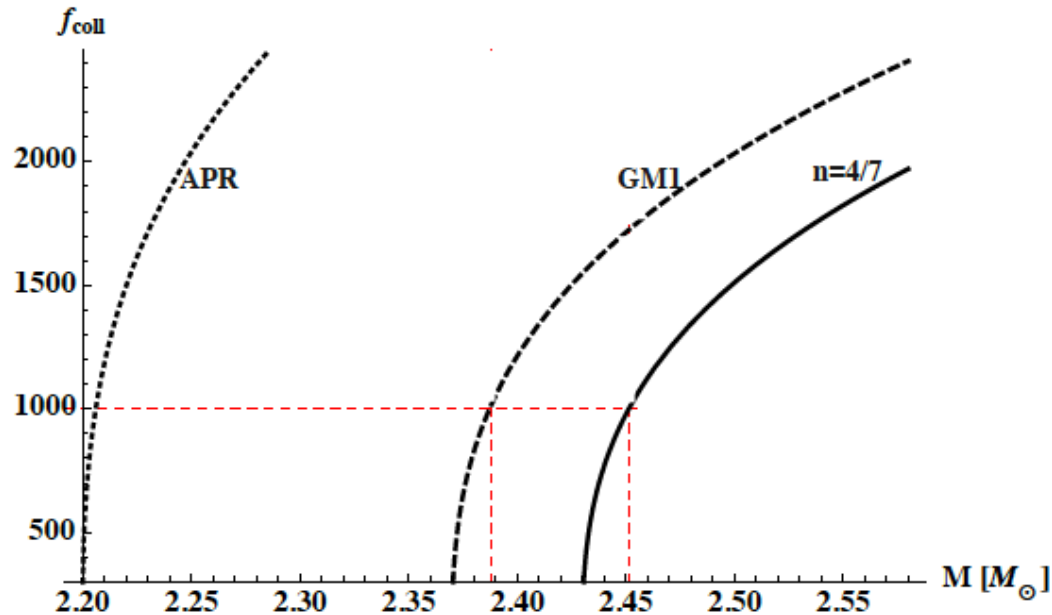


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Fast (uniform) rotation  $\rightarrow$  sustains supramassive PMNS against collapse  
Falcke & Rezzolla (2014), Lasky et al. (2014)





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**2) Event Rate: How many BNS mergers?**

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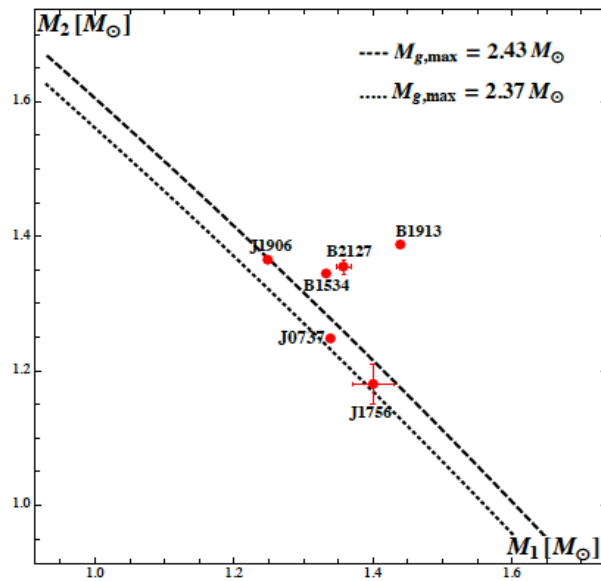
- a)* distance of expected signals  $\leq 35$  Mpc
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- c)* fraction of stable NSs depends on EOS

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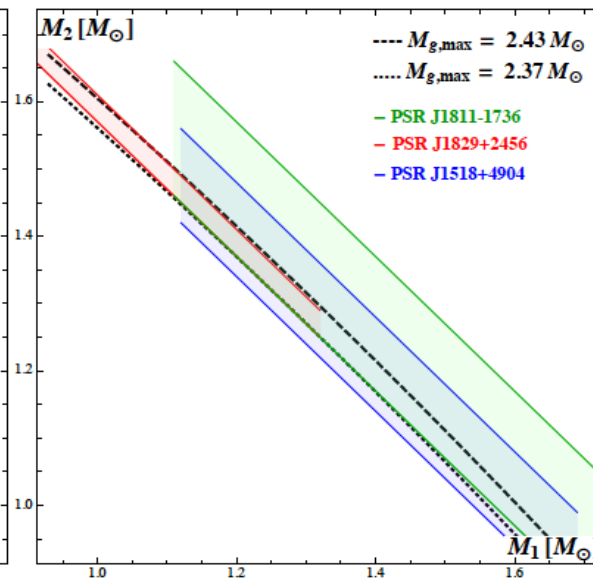
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Dall'Osso et al. (2014)

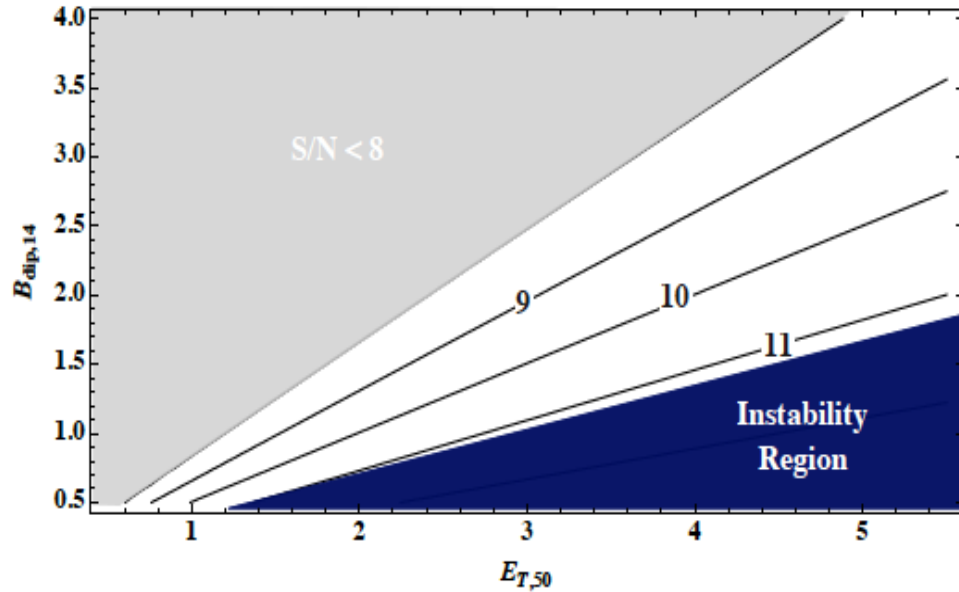
data from Kiziltan et al. (2013)



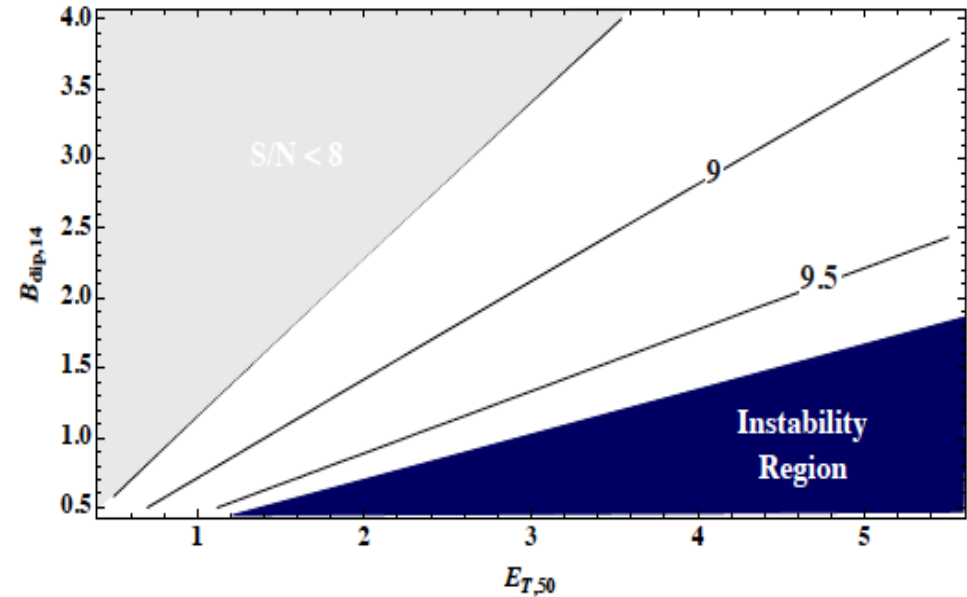
$$M_{rest} = M_g + 0.075 M_g^2$$

Timmes et al. (1996)

# Massive “Magnetars” from BNS mergers



Stable PMNS  
 $D \sim 35$  Mpc



Hypermassive PMNS  
 $D \sim 16$  Mpc

Dall'Osso et al. (2014)

# Open questions

Dedicated observational strategies to reveal the putative signals (ongoing)

Generalization of results to interior field configuration of “arbitrary” shape and for different EOS's (next step)

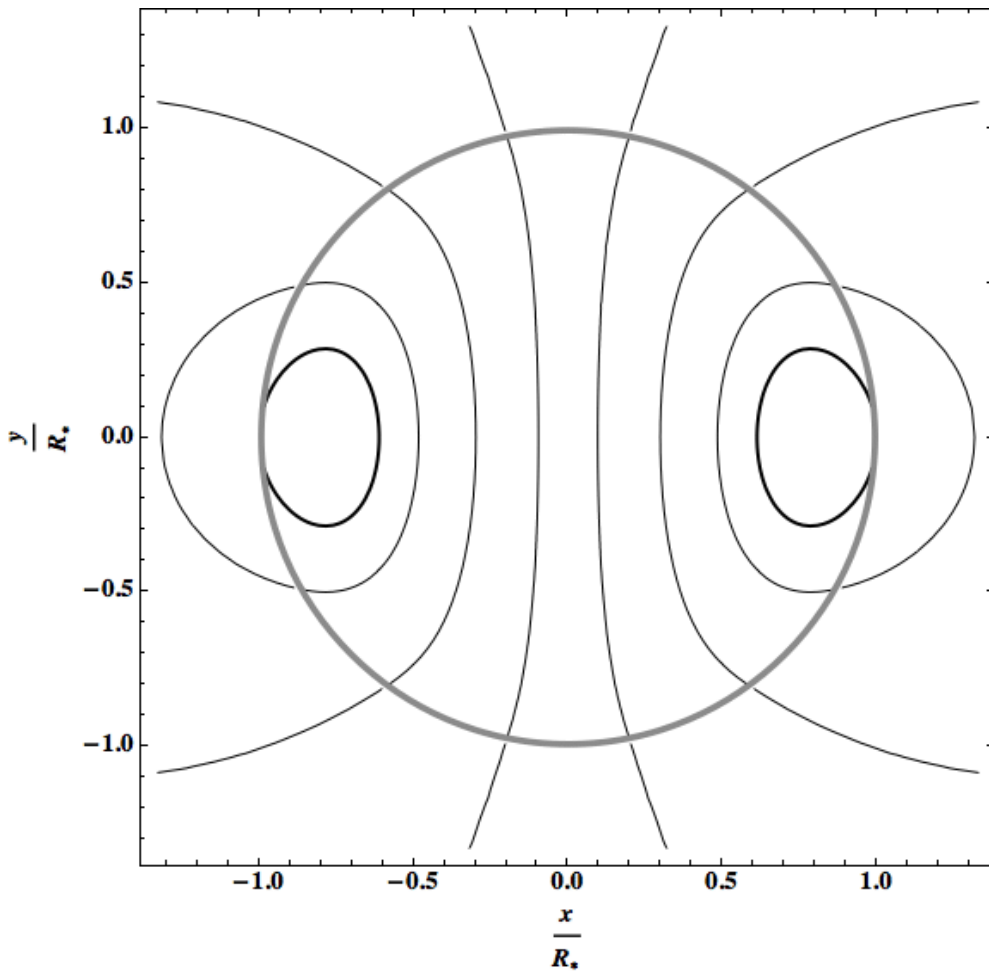
Calculation of microphysical effects beyond current approximation, e.g. stability of magnetic fields, the effect of internal viscosity, etc.

Understanding the fate of angular momentum and the initial spin rate of the PMNS, particularly important for “unstable” PMNS

...and of course determining the **actual event rates. GW signals!?**

# **Extra Slides**

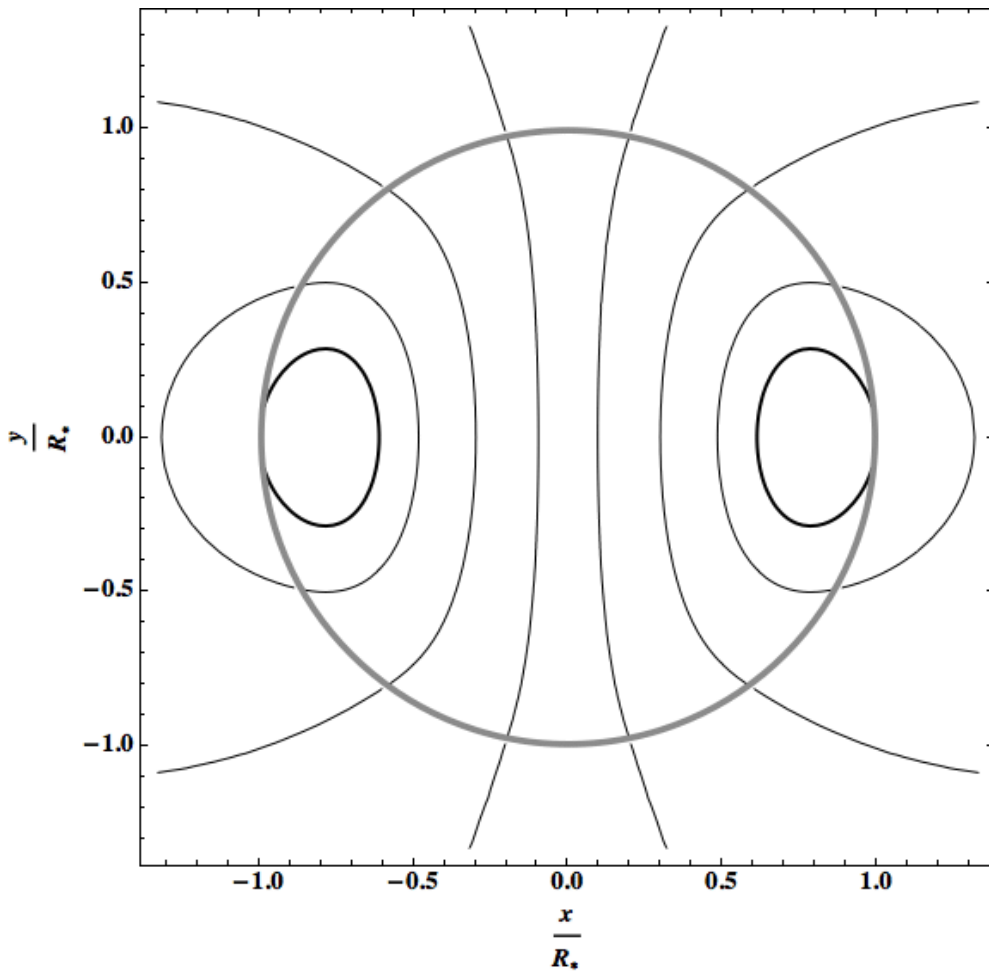
# Twisted-Torus Configuration



$$(1/\mu_0)(\nabla \times \mathbf{B}) \times \mathbf{B} = \nabla \delta p + \delta \rho \nabla \Phi,$$

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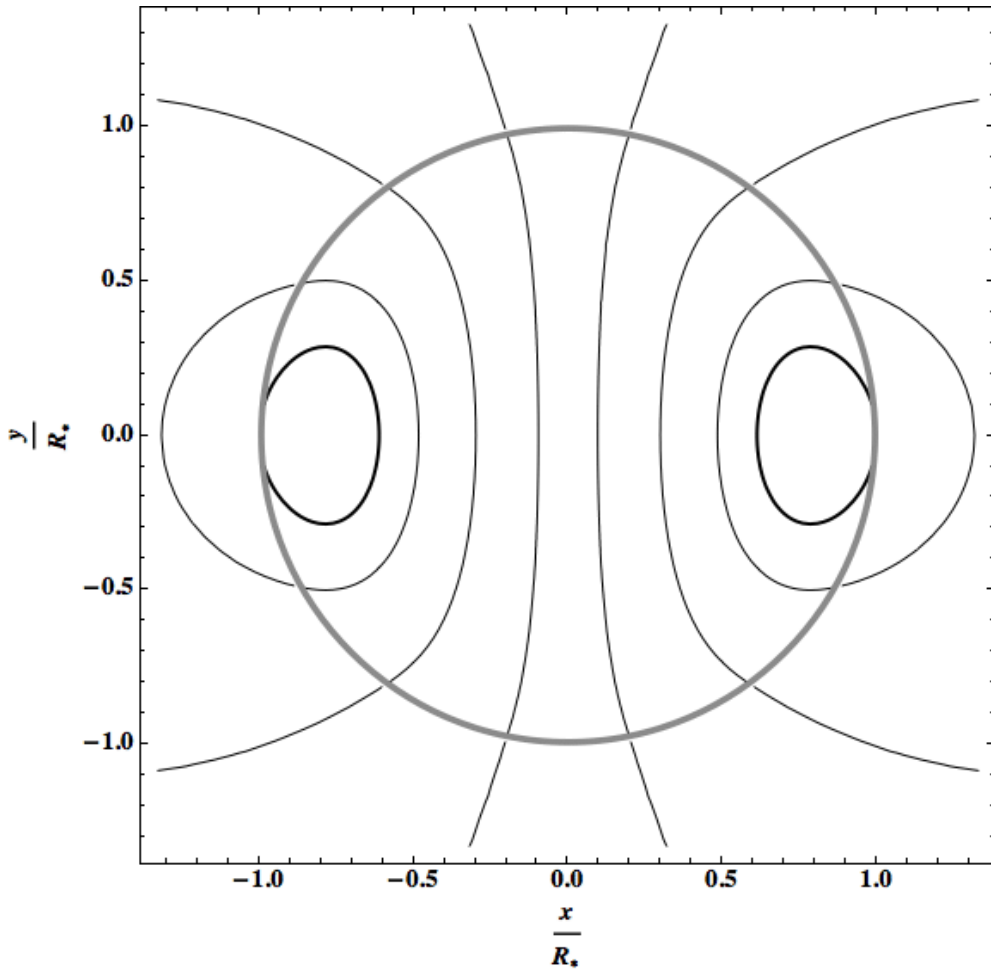


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# Twisted-Torus Configuration

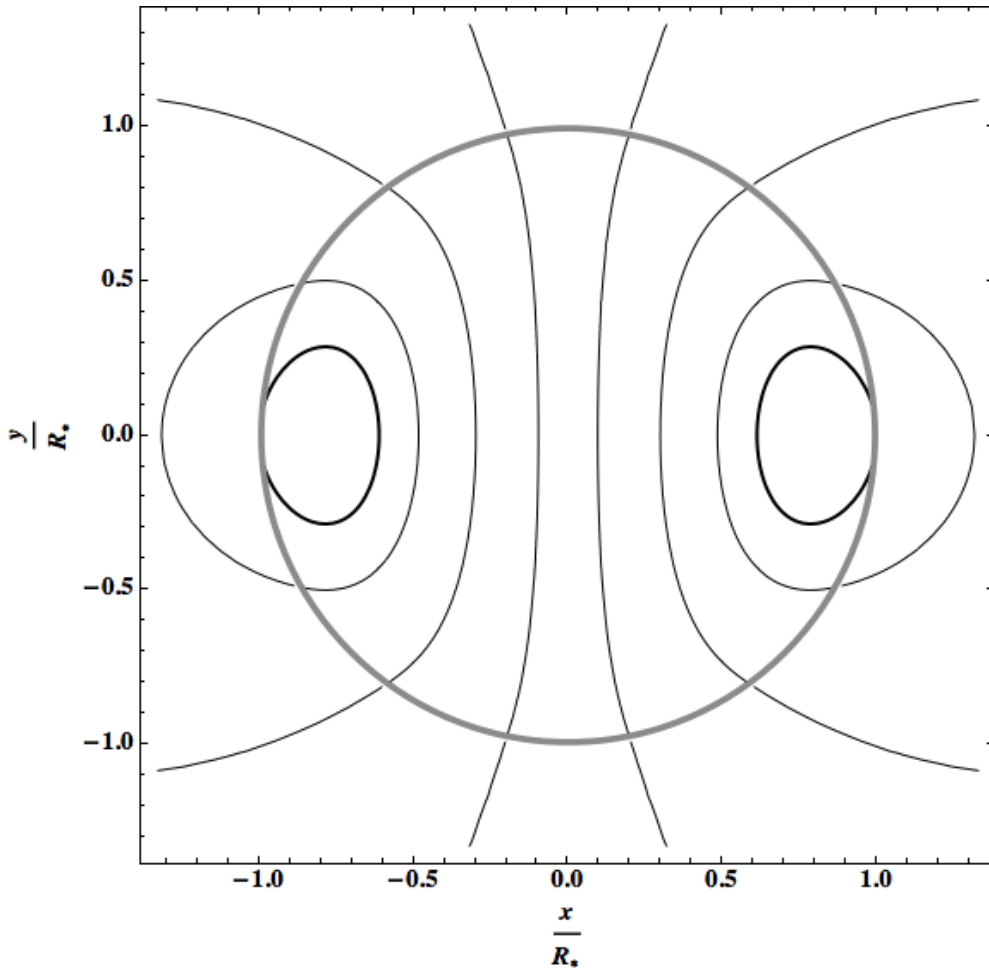


$$(1/\mu_0)(\nabla \times \mathbf{B}) \times \mathbf{B} = \nabla \delta p + \delta \rho \nabla \Phi,$$

$$\mathbf{B} = B_0[\eta_p \nabla \alpha(r, \theta) \times \nabla \phi + \eta_t \beta(\alpha) \nabla \phi],$$

$$-\frac{B_0^2}{\mu_0 r^2 \sin^2 \theta} (\eta_p^2 \nabla \alpha \hat{\Delta} \alpha + \eta_t^2 \beta \nabla \beta) = \nabla \delta p + \delta \rho \nabla \Phi,$$

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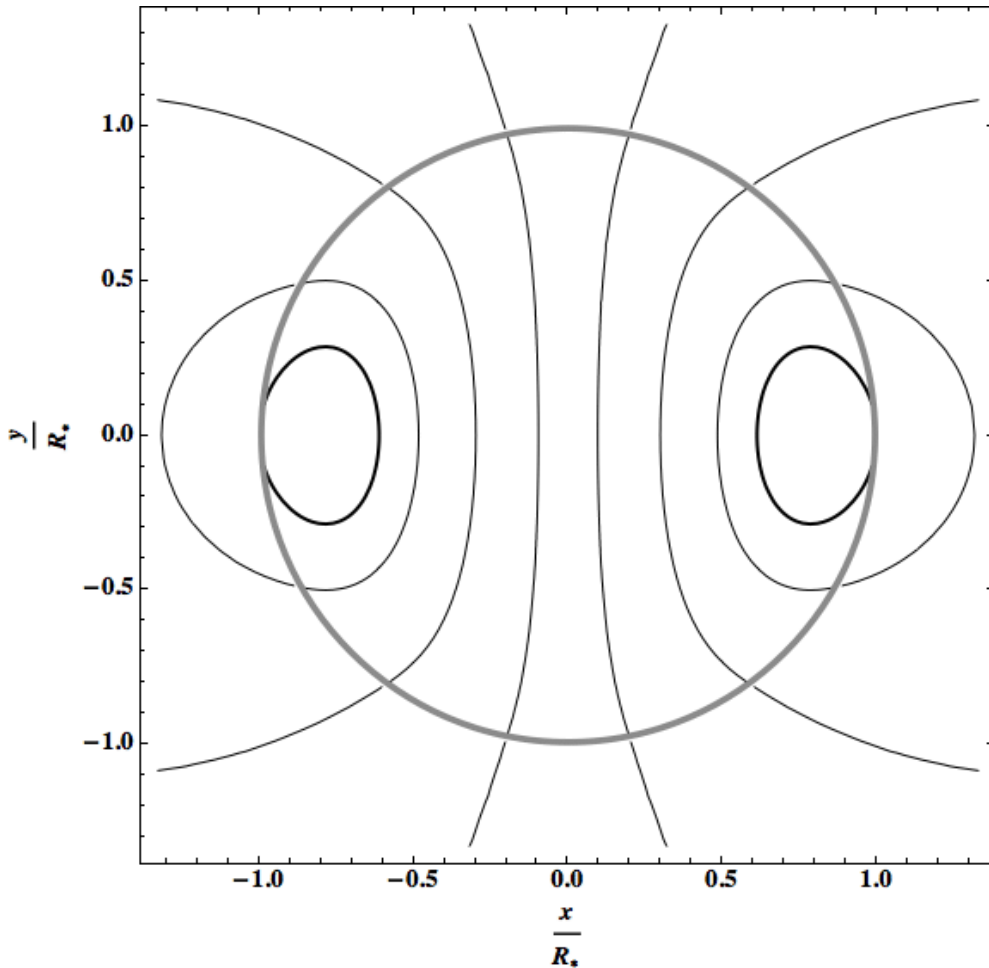
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$$\hat{\Delta} = \frac{\partial^2}{\partial r^2} + \frac{\sin \theta}{r^2} \frac{\partial}{\partial \theta} \left( \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \right)$$

# Twisted-Torus Configuration



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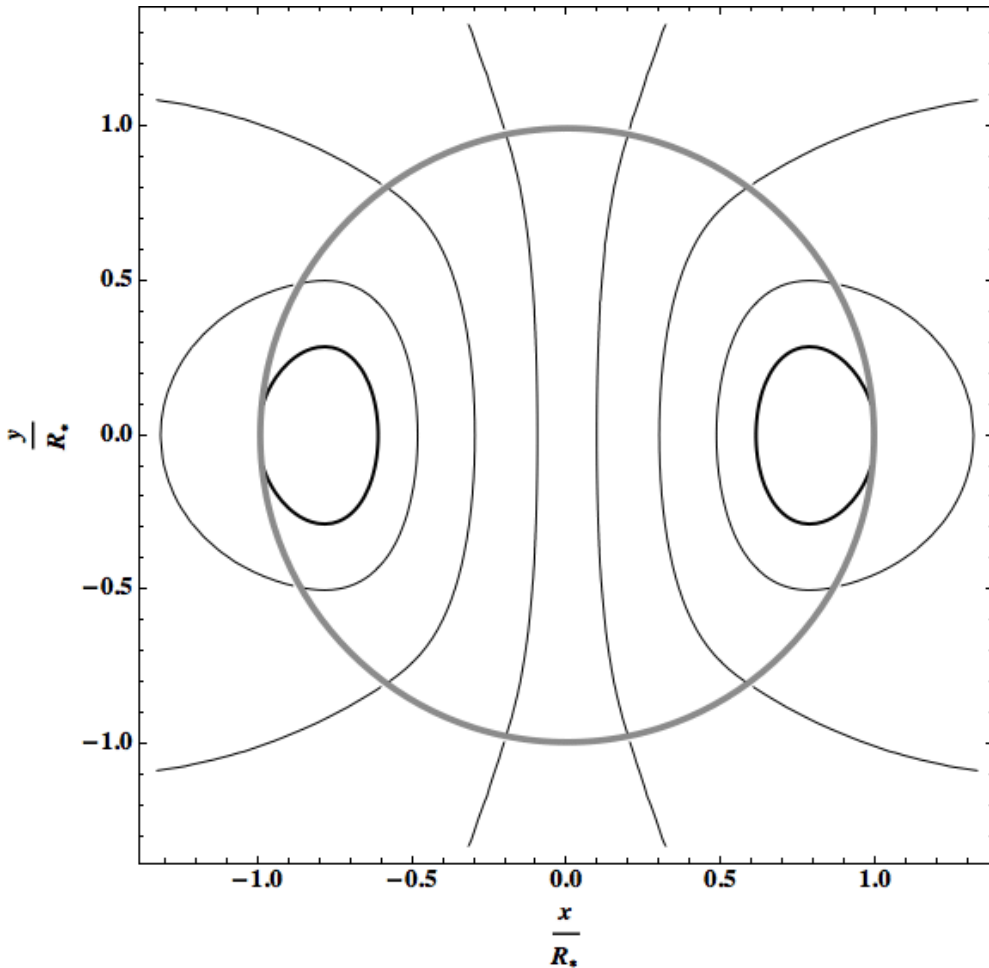
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$$\delta p = \frac{-B_0^2 \alpha}{\mu_0 r^2 \sin^2 \theta} \left[ \eta_p^2 \hat{\Delta} \alpha + 2\eta_t^2 \left( \frac{\alpha^3}{3} - \frac{3\alpha^2}{2} + 3\alpha - \ln \alpha + C_0 \right) \right]$$

# Twisted-Torus Configuration



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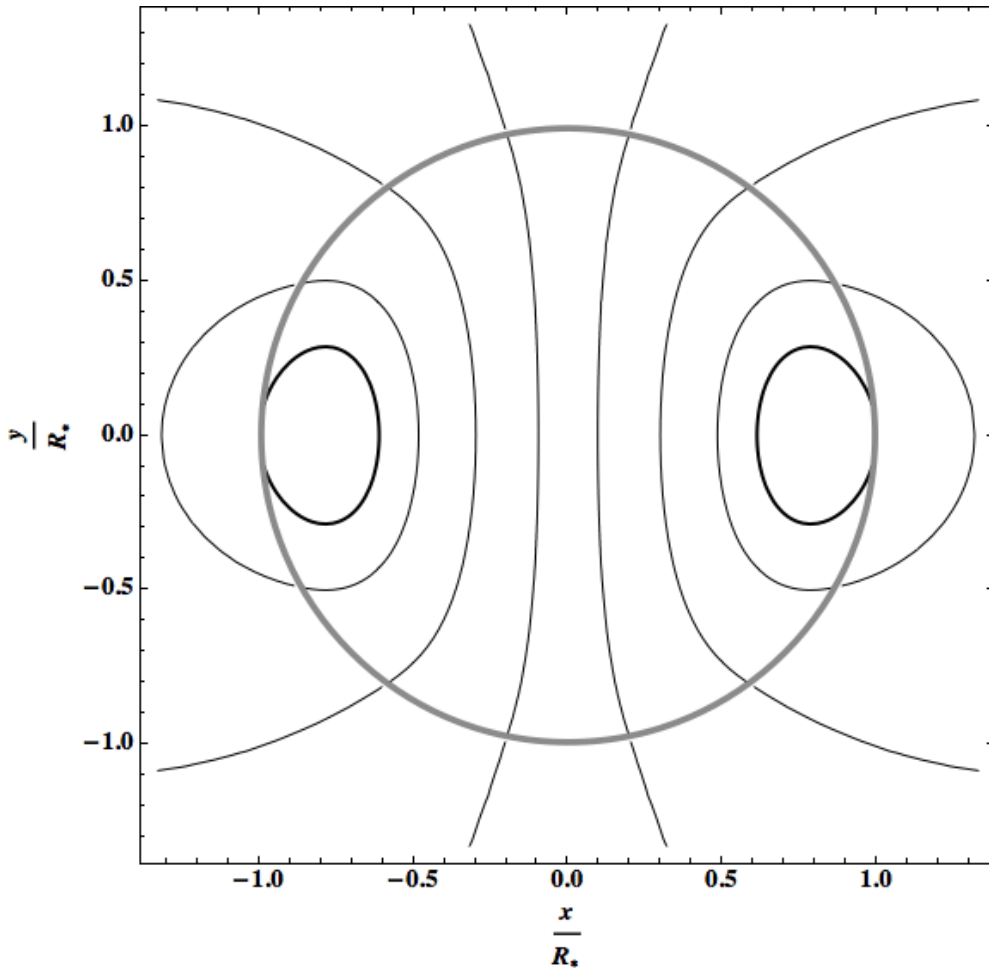
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$$\delta \rho \frac{d\Phi}{dr} = \frac{B_0^2}{\mu_0} \left[ \eta_p^2 \alpha \frac{\partial}{\partial r} \left( \frac{\hat{\Delta} \alpha}{r^2 \sin^2 \theta} \right) + 2\eta_t^2 \left( \frac{\alpha^3}{3} - \frac{3\alpha^2}{2} + 3\alpha - \ln \alpha + C_0 \right) \frac{\partial}{\partial r} \left( \frac{\alpha}{r^2 \sin^2 \theta} \right) \right]$$

# Twisted-Torus Configuration



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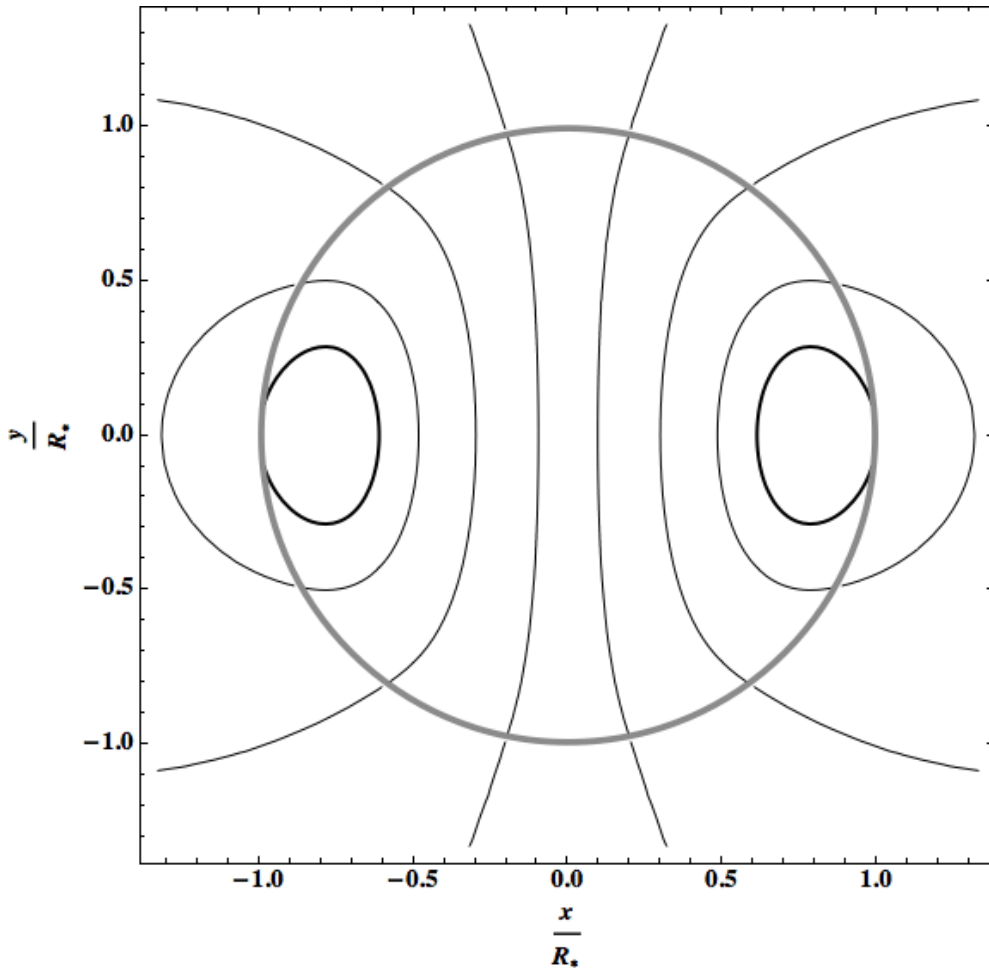
$$-\frac{B_0^2}{\mu_0 r^2 \sin^2 \theta} (\eta_p^2 \nabla \alpha \hat{\Delta} \alpha + \eta_t^2 \beta \nabla \beta) = \nabla \delta p + \delta \rho \nabla \Phi,$$

$$I_{jk} = R_*^5 \int_V dV [\rho(r) + \delta \rho(r, \theta)] (r^2 \delta_{jk} - x_j x_k).$$

$$\epsilon = \frac{I_{zz} - I_{xx}}{I_0}$$

$$\delta \rho \frac{d\Phi}{dr} = \frac{B_0^2}{\mu_0} \left[ \eta_p^2 \alpha \frac{\partial}{\partial r} \left( \frac{\hat{\Delta} \alpha}{r^2 \sin^2 \theta} \right) + 2\eta_t^2 \left( \frac{\alpha^3}{3} - \frac{3\alpha^2}{2} + 3\alpha - \ln \alpha + C_0 \right) \frac{\partial}{\partial r} \left( \frac{\alpha}{r^2 \sin^2 \theta} \right) \right]$$

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$$\epsilon = \pi R_*^5 I_0^{-1} \int_V dr d\theta \delta \rho(r, \theta) r^4 \sin \theta (1 - 3 \cos^2 \theta)$$

$$\delta \rho \frac{d\Phi}{dr} = \frac{B_0^2}{\mu_0} \left[ \eta_p^2 \alpha \frac{\partial}{\partial r} \left( \frac{\hat{\Delta} \alpha}{r^2 \sin^2 \theta} \right) + 2\eta_t^2 \left( \frac{\alpha^3}{3} - \frac{3\alpha^2}{2} + 3\alpha - \ln \alpha + C_0 \right) \frac{\partial}{\partial r} \left( \frac{\alpha}{r^2 \sin^2 \theta} \right) \right]$$